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**Interim Report**

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**Mutant invasions and adaptive dynamics  
in variable environments**

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1 **Mutant invasions and adaptive dynamics in variable**

2 **environments**

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17 ***Abstract***

18 The evolution of natural organisms is ultimately driven by the invasion and possible  
19 fixation of mutant alleles. The invasion process is highly stochastic, however, and the  
20 probability of success is generally low, even for advantageous alleles. Additionally, all  
21 organisms live in a stochastic environment, which may have a large influence on what  
22 alleles are favourable, but also contributes to the uncertainty of the invasion process. We  
23 calculate the invasion probability of a beneficial mutant allele in a monomorphic, large  
24 population subject to stochastic environmental fluctuations, taking into account density  
25 and frequency dependent selection, stochastic population dynamics and temporal  
26 autocorrelation of the environment. We treat both discrete and continuous time  
27 population dynamics, and allow for overlapping generations in the continuous time case.  
28 The results can be generalized to diploid, sexually reproducing organisms embedded in  
29 communities of interacting species. We further use these results to derive an extended  
30 canonical equation of adaptive dynamics, predicting the rate of evolutionary change of a  
31 heritable trait on long evolutionary time scales.

32 ***Introduction***

33 Although the ecological importance and basic principles of adaptation to a variable  
34 environment have been long known, the corresponding genetic processes are not yet  
35 sufficiently understood. Ultimately, evolution is dependent on the fate of mutant alleles,  
36 and during the first generations after the appearance of a new variety its success is to a

37 large extent dependent on chance events and the probability of extinction is high. A large  
38 body of theory (nicely reviewed by Patwa & Wahl (2008)) treats the probability that an  
39 advantageous mutant survives the first crucial generations and becomes sufficiently  
40 abundant so that the risk of stochastic extinction can be ignored. This has in the literature  
41 been called the probability of ‘survival’, ‘establishment’, ‘fixation’ or ‘invasion’,  
42 depending on the context. We will here use the term ‘invasion’. In many cases invasion  
43 implies fixation, but not necessarily so if fitness is frequency dependent, such that a  
44 polymorphism is possible.

45 Starting with the simpler case of a constant environment, Haldane (1927) famously stated  
46 that the invasion probability of a mutant allele equals  $2s$ , where  $s$  is the relative fitness  
47 advantage of the invading allele (Haldane assumed a constant, large population size,  
48 Poisson distribution of offspring and a small  $s$ ). Later, Ewens (1969) and Eshel (1981)  
49 (see also Athreya (1992)) generalized Haldane’s result to arbitrary offspring distributions.  
50 They found the invasion probability to be approximately equal to  $2s/\sigma^2$ , where  $\sigma^2$  is the  
51 variance in the number of offspring from a single individual, i.e. a measure of the  
52 strength of genetic drift (or demographic stochasticity). For example the Poisson  
53 distribution has a variance equal to its mean, which by assumption is equal to  $1+s$  here.  
54 Thus, Ewens’ and Eshel’s approximation agrees with Haldane’s result since  $s$  is assumed  
55 to be small.

56 Taking variable survival and/or reproduction rate into account is inherently difficult in  
57 the general case. The case of a variable fitness advantage  $s$  but constant population size  $N$

58 has been studied several times (e.g. Kimura 1954, Jensen 1973, Karlin & Levikson 1974,  
59 Takahata et al. 1975). Alternatively, a branching process approach can be used, which  
60 usually requires the assumption of an infinite resident population size. Smith &  
61 Wilkinson (1969) showed by this approach that an invading mutant will go extinct with  
62 certainty if  $\mathbb{E}(\ln(m_t)) < 0$ , where  $m_t$  is the time-dependent average number of offspring per  
63 individual and  $\mathbb{E}(\cdot)$  denotes the long term, stationary, mean (Dempster 1955 fore-  
64 shadowed this result, see also Gillespie 1973). It is assumed that each  $m_t$  is chosen  
65 independently from a fixed distribution – a so-called white noise environment. Later,  
66 Athreya & Karlin (1971) generalized this result to autocorrelated environments, and  
67 Karlin & Lieberman (1974) to diploid populations. Together, these results underline the  
68 importance of mean log growth rate for adaptations to variable environments, a  
69 fundamental result in bet-hedging theory (e.g. Cohen 1966, Seger & Brockman 1987). In  
70 a recent paper, Peischl & Kirkpatrick (2012) used novel analytical techniques to calculate  
71 the probability of invasion, given small fluctuations of  $s$ . They show that the invasion  
72 probability is proportional to a weighted time-average of  $s$ , with more weight on points in  
73 time with low mutant abundance.

74 If the invading mutant has a fixed fitness advantage relative to the resident type, then the  
75 mutant growth rate will vary over time just like that of the resident population. This  
76 assumption has been used in a number of studies. Ewens (1967) showed that the  
77 probability of establishment in a cyclic population equals  $2s \frac{n_H}{n(0)}$  (again assuming a

78 Poisson distribution of offspring and a small  $s$ ), where  $n_H$  is the harmonic mean  
79 population size and  $n(0)$  is the resident population size at the time when the mutant first  
80 appears. This shows that the invasion of a mutant type is less likely if the amplitude of  
81 the population cycle is large (assuming a fixed arithmetic mean), since the harmonic  
82 mean is sensitive to variation, as opposed to the arithmetic mean. It can also be shown  
83 that invasion is more likely in a growing population than in a declining population  
84 (Ewens 1967, Kimura & Ohta 1974, Otto & Whitlock 1997). The results by Ewens  
85 (1967) and Otto & Whitlock (1997) for cyclic populations were later generalised to  
86 arbitrary offspring distributions by Pollak (2000), who among other things confirmed that  
87 the probability of invasion in a cyclic population is proportional to the harmonic mean  
88 population size divided by the population size at mutant introduction.

89 The more general case of both a variable strength of selection and a variable resident  
90 population size has been treated recently by Waxman (2011), Uecker & Hermisson  
91 (2011). In both studies, quite general, but rather implicit, expressions for the invasion  
92 probability are derived. Uecker & Hermisson further analyze simplifying cases such as a  
93 deterministically growing population or a periodic (sinusoidal) environment.

94 Lastly, we would like to highlight a rarely cited result by Hill (1972) who, somewhat  
95 offhandedly, derived the expression

96 
$$P = \frac{1 - e^{-2n_c \bar{s} q}}{1 - e^{-2n_c \bar{s}}}, \quad (1)$$

97 where  $P$  is the probability of mutant invasion,  $n_e = n_H$  is again the harmonic mean  
98 population size,  $\bar{s}$  is the arithmetic mean selective advantage and  $q$  is the initial  
99 proportion of the mutant type. We will return to this result, and its assumptions, in later  
100 sections.

101 We here generalize several of the aforementioned results to the case of arbitrary ergodic  
102 population dynamics, subject to ergodic environmental fluctuations. We calculate the  
103 invasion probability of a mutant of small phenotypic effect in a large resident population.  
104 Mutant fitness, and in particular its selective advantage  $s$ , depends on the resident  
105 population size as well as the environmental fluctuations and may in some circumstances  
106 be negative as long as the long term mean ( $\bar{s}$ ) is positive. Solutions are given for both  
107 discrete time and continuous time dynamics. The continuous time case allows for  
108 overlapping generations and is a particularly suitable model for unicellular organisms that  
109 reproduce through fission, such as bacteria or protozoa.

### 110 ***Model description, basic assumptions***

111 We consider the invasion of a mutant type in a monomorphic resident population of  
112 asexually reproducing individuals, under the assumptions that *i*) all individuals are  
113 equivalent, i.e. there is no age-, stage- or spatial structure, *ii*) the resident population size  
114 is large enough that the growth of an invading mutant is independent of its own density,  
115 at least until the mutant abundance is large enough that the risk of stochastic extinction is  
116 negligible, and *iii*) the mutation is of small effect, such that the mutant type is

117 ecologically close to the resident type, i.e. it has in all possible environmental  
118 circumstances a per capita growth rate close to that of the resident.

## 119 Concepts and notation

120 Since we will move back and forth between the established conceptual frameworks of  
121 stochastic population dynamics, population genetics and long term evolution, a couple of  
122 concepts may have different meanings to readers with different background.

123 Firstly, the *environment* of an invading mutant type consists of two basic components –  
124 the *external environment* and the *feedback environment*. We think of the external  
125 environment as a stochastic, ergodic process, which affects the survival and reproductive  
126 success of all individuals of the same type in the same way, such as stochastic weather  
127 fluctuations or a variable resource abundance. *Ergodic* means that irrespective of initial  
128 conditions, the environment will in the long term visit its full stationary distribution. The  
129 external environment is in itself not affected by the state of the focal population, in  
130 contrast to the feedback environment, which by definition depends on the current state of  
131 the focal population and possible interacting populations (Metz et al 1992, Mylius &  
132 Diekmann 1995, Heino et al. 1998). In the simplest of cases the feedback environment is  
133 population size and the external environment is a single parameter, such as temperature.  
134 Our analysis is staged in this simplified scenario but it is straightforward to generalize to  
135 the multidimensional case (see below).

136 Secondly, *fitness* can be understood either as long term fitness, i.e. the long term average  
137 per capita growth rate of any given clone, or as the instantaneous per capita growth rate at  
138 any given moment. We use the qualifications *mean fitness* and *instantaneous fitness* to  
139 denote the two concepts, respectively (more precise definitions follow).

140 Finally, we use  $\mathbb{E}[z(t)]$ ,  $\mathbb{V}[z(t)]$  and  $\mathbb{C}[z(t),w(t)]$  to denote the mean, variance and  
141 covariance, respectively, of the stochastic process(es)  $z(t)$  (and  $w(t)$ ). If nothing else is  
142 specified, the stationary mean, variance and covariance, respectively, are intended. For  
143 brevity, we will sometimes use  $\bar{z}$  to denote the mean.

#### 144 Continuous time model

145 We start with the continuous time case – assuming individuals reproduce and die  
146 according to a time-inhomogeneous birth and death process. More formally, we assume  
147 that a resident type individual has a birth rate,  $b(n(t), \varepsilon(t))$ , and death rate  $d(n(t), \varepsilon(t))$ ,  
148 where  $n(t)$  is the resident population size and  $\varepsilon(t)$  is an environmental process. It is  
149 assumed that  $\varepsilon(t)$  is an ergodic, stochastic process continuous in time. The instantaneous  
150 fitness, i.e. the per capita growth rate,  $f$ , is given by the difference between birth and  
151 death rate,

$$152 \quad f(n(t), \varepsilon(t)) = b(n(t), \varepsilon(t)) - d(n(t), \varepsilon(t)). \quad (2)$$

153 We denote the total dynamic environment determining the instantaneous fitness  $E(t)$ . In  
154 the formalism here,  $E(t) = \{n(t), \varepsilon(t)\}$  and the growth, birth and death rates can be written

155 
$$f(E(t)) = b(E(t)) - d(E(t)) \tag{3}$$

156 We assume that  $E(t)$  is ergodic, which should be a realistic assumption for many  
157 scenarios, albeit excluding long-term environmental trends or a steadily growing or  
158 declining population. Note that autocorrelation of the environmental process  $\varepsilon(t)$  is  
159 allowed, as long as it declines to zero at large time lags. More precisely, the total  
160 environment  $E(t)$  should explore its full stationary distribution much faster than the time  
161 scale of a mutant invasion ( $1/\bar{s}$ , see below). It should also be noted that technically  
162 speaking the population process is not ergodic since  $n = 0$  is an absorbing state. However,  
163 in the large population limit considered here, this is of minor importance.

164 Given the growth function above, it is straightforward to express the resulting dynamics  
165 of the resident population. Since we assume population size  $n$  to be large enough that  
166 demographic stochasticity can be ignored, the resident population dynamics are given by

167 
$$\frac{dn}{dt} = f(E(t))n(t) \tag{4}$$

168 We assume a single mutant individual appears in the population at  $t = 0$ . The mutant  
169 birth, death and per capita growth rates are denoted  $\tilde{b}(E(t))$ ,  $\tilde{d}(E(t))$  and  $\tilde{f}(E(t))$   
170 respectively. The instantaneous mutant fitness advantage is written

$$171 \quad s(E(t)) = \tilde{f}(E(t)) - f(E(t)). \quad (5)$$

172 Note that  $E(t)$  is still the environment given by the population dynamics of the resident  
173 population (and the external environment). A mutant type may have a fixed fitness  
174 advantage ( $s$ ), but can also differ in its density dependence, its sensitivity to fluctuations  
175 of the external environment, or all of the above.  $s(E(t))$  can in the general case change  
176 sign depending on the state of the environment  $E(t)$ , but we assume its long-term  
177 (stationary) mean,  $\bar{s}$ , is positive. In other words, the mutant type may be at a  
178 disadvantage for shorter periods of time, as long as it is advantageous on average.

### 179 Discrete time model

180 For the discrete time case we assume non-overlapping generations. Each individual  
181 (independently) gives birth to a geometrically distributed number of offspring, with the  
182 mean number of offspring determined by the individual's instantaneous fitness. The  
183 probability of  $k$  offspring is

$$184 \quad \Pr(k) = (1 - p)^k p, \quad (6)$$

185 where  $p = 1/(1+\lambda)$  and  $\lambda$  is the mean number of offspring. The variance in offspring  
186 number is  $\lambda(\lambda + 1)$ , which can be compared to the commonly used Poisson distribution,  
187 which has a variance equal to its mean,  $\lambda$ . A mechanistic motivation for the geometric  
188 distribution arises if an individual makes repeated reproduction attempts, each with the  
189 same probability of success, but stops at the first failure. From a more pragmatic point of  
190 view, however, there is clearly no natural population where individual reproductive  
191 success exactly follows a geometric or Poisson distribution. The geometric distribution is  
192 used here for mathematical convenience, in lack of a more general theory for all, or at  
193 least a family of distributions.

194 In discrete time we define the instantaneous fitness function  $f$  as the natural logarithm of  
195 the per-capita growth rate ( $\lambda$ ), such that the mean number of surviving offspring of an  
196 individual of the resident type is given by  $e^{f(n(t), \varepsilon(t))} = e^{f(E(t))}$ , where  $\varepsilon(t)$  here is a discrete  
197 time process, but with otherwise the same properties as in the continuous time case  
198 above. The dynamics of a large population of resident type individuals is thus

$$199 \quad n(t+1) = e^{f(E(t))} n(t) . \tag{7}$$

## 200 ***Mutant invasion***

201 We here derive the main result – the probability of invasion of a mutant type, starting as a  
202 single individual at time  $t = 0$ . Invasion does not necessarily imply fixation. If  
203 coexistence of the mutant and resident types is possible, we assume the equilibrium

204 mutant abundance is large, such that the invasion process can safely be analyzed under  
 205 the assumption that mutant abundance has no effect on mutant fitness. More precisely, we  
 206 assume there is a population size  $n_i$  of the mutant type at which invasion can be  
 207 considered certain but that at the same time  $n_i \ll n$ , where  $n$  is the equilibrium resident  
 208 population size. If the probability that a mutant population starting with a single  
 209 individual invades is equal to  $P$ , then the probability that a population of  $n_i$  mutants goes  
 210 extinct is approximately given by  $(1 - P)^{n_i} \approx e^{-n_i P}$  as long as  $P$  is small. A requirement is  
 211 thus that  $e^{-n_i P}$  is close to zero, i.e. that  $n_i P$  is large ( $n_i P > 5$  gives an error less than 1%).  
 212 If, as we will show,  $P$  is the size of  $\bar{s}$  we can express the necessary requirement that  
 213  $n\bar{s} \gg 1$  for our analysis to hold.

#### 214 Continuous time

215 As a starting point, we use a result by Kendall (1948), which states

$$216 \quad P_E = \frac{1}{1 + I_E}, \quad (8a)$$

217 where  $P_E$  is the ultimate survival probability of a time-dependent birth-and-death process  
 218 and

$$219 \quad I_E = \int_0^\infty \tilde{d}(E(t)) e^{-\int_0^t \tilde{f}(E(\tau)) d\tau} dt. \quad (8b)$$

220 A heuristic interpretation of equation (8b) is a weighted total death rate, with most weight  
 221 on periods, usually at low  $t$ -values, with low numbers of mutants (the exponential factor  
 222 can be interpreted as  $1/(\text{expected mutant population size at time } t)$ ). As mentioned in the  
 223 introduction, a similar weighting was found by Peischl and Kirkpatrick (2012).

224 The environment  $E(t)$  is in the general case stochastic and unpredictable. The necessary  
 225 interpretation of  $P_E$  (eq. 8a) is thus the *conditioned* survival probability (Waxman 2011),  
 226 conditioned on the future environment  $E(t)$ ,  $t \geq 0$ , which is the reason for the subscript  $E$ .

227 The unconditioned probability of invasion is given by the mean  $P_E$ , and we here calculate  
 228 the mean probability  $P_0$ ,

$$229 \quad P_0 = \mathbb{E}[P_E | E(0)], \tag{9}$$

230 averaged across all possible future developments of environmental states, but still  
 231 conditioned on initial conditions  $E(0)$ . In particular, we seek the linear dependence of  $P_0$   
 232 on the mean fitness advantage  $\bar{s}$  as  $\bar{s}$  becomes small, i.e. we seek the limit

$$233 \quad \lim_{\bar{s} \rightarrow 0} \frac{P_0}{\bar{s}} = \lim_{\bar{s} \rightarrow 0} \mathbb{E} \left[ \frac{P_E}{\bar{s}} | E(0) \right] = \lim_{\bar{s} \rightarrow 0} \mathbb{E} \left[ \frac{1}{\bar{s} + \bar{s} I_E} | E(0) \right]. \tag{10}$$

234 In Appendix A we show that

235 
$$\lim_{\bar{s} \rightarrow 0} \bar{s} I_E = n(0) \mathbb{E} \left[ \frac{d(E(t))}{n(t)} \right] \quad (11)$$

236 for almost all possible future environments  $E(t)$ ,  $t \geq 0$  (the exceptions have probability  
 237 zero).  $n(0)$  is the resident population size at the time of mutant arrival, but all other  
 238 dependencies on initial conditions average out. Inserting equation (11) into equation (10)  
 239 gives (see Appendix A for details)

240 
$$\lim_{\bar{s} \rightarrow 0} \frac{P_0}{\bar{s}} = \frac{1}{n(0) \mathbb{E} \left[ \frac{d(E(t))}{n(t)} \right]} = \frac{1}{n(0)(d/n)}, \quad (12)$$

241 and we can finally express the approximate invasion probability as

242 
$$P_0 \approx \frac{\bar{s}}{n(0)(d/n)} = 2 \frac{\bar{s}}{b} \frac{n_e}{n(0)}, \quad (13a)$$

243 where we define the effective population size  $n_e$  as

244 
$$n_e = \frac{\bar{d}}{2(d/n)} = \frac{\bar{b}}{2(b/n)} = \frac{\overline{(b+d)}}{2(b+d)/n}. \quad (13b)$$

245 The identities  $\bar{b} = \bar{d}$  and  $\overline{(d/n)} = \overline{(b/n)} = \frac{1}{2}\overline{(b+d)/n}$  follow from the ergodicity  
246 assumption of  $n(t)$ . More precisely, they follow from the assumptions that  $\ln(n(t))$  and  
247  $1/n(t)$  have a long term mean growth rate of zero.

248 The definition of effective population size (eq. 13b) is somewhat arbitrary. Otto &  
249 Whitlock (1997) suggest defining  $n_e$  such that  $P_0 = 2\bar{s}n_e / n(0)$  (the “fixation effective  
250 population size”), which in our case implies setting  $n_e = 1/(2\overline{b/n})$ . However, our  
251 proposed definition of effective population size (eq. 13b) has the appealing properties that  
252 *i)* it is unitless – it does not depend on the chosen time unit, *ii)* it simplifies to  $n_e = n/2$  in  
253 cases when  $n$  is constant, *iii)* it can be interpreted as half the weighted harmonic mean  
254 population size, weighted by the total per capita event rate  $(b + d)$ , and is thus congruent  
255 with the discrete time case below. A possible disadvantage with our definition is that the  
256 average fitness advantage,  $\bar{s}$ , must be standardized with the mean birth rate,  $\bar{b}$ . On the  
257 other hand, the unitless ratio  $\bar{s}/\bar{b}$  (Eq. 13a) can be interpreted as a standardized selection  
258 coefficient, measured on the time scale of the average generation time (in the  
259 deterministic case, with a constant population size, generation time equals  $1/d = 1/b$ ).  
260 Irrespective of the preferred definition of effective population size, equation (13a) is  
261 directly comparable to several previous results in discrete time (e.g. Ewens 1967, Otto &  
262 Whitlock 1997, Pollack 2000).

263 The approximation in equation (13a) is valid for small  $s$ , i.e. not only is  $\bar{s}$  small, but also  
264 its fluctuations. The mutant type can thus not be inherently different from the resident

265 type – its instantaneous fitness must for all environmental states be close to that of the  
266 resident. The only realistic interpretation is a mutation of small phenotypic effect. We  
267 further investigate the applicability of this result in the Model Examples section below  
268 and in Appendix C (online supplement).

### 269 The discrete time case

270 Using the assumption of geometrically distributed offspring, the ultimate survival  
271 probability of a mutant strategy appearing at  $t = 0$  can be expressed exactly as (Haccou et  
272 al. 2005, Box 5.5):

$$273 \quad P_E = \frac{1}{1 + I_E} \quad (14a)$$

274 where

$$275 \quad I_E = \sum_{t=0}^{\infty} e^{-\sum_{\tau=0}^t \tilde{f}(E(\tau))} . \quad (14b)$$

276 The striking similarity between equations (14a,b) and the continuous time version  
277 equations (8a,b) makes it possible to carry out almost exactly the same derivation as  
278 above, only exchanging integrals with sums and setting the death rates  $d$  and  $\tilde{d}$  to 1. Due  
279 to the great similarity of the calculations we refrain from presenting the discrete time  
280 derivation here, and instead present the major results:

281  $P_0 \approx 2\bar{s} \frac{n_e}{n(0)},$  (15a)

282 where

283  $n_e = \frac{n_H}{2},$  (15b)

284 and, just like above,  $P_0$  is the probability of invasion conditioned on initial conditions  
285  $E(0)$ ,  $n(0)$  is the resident population size at the time of mutant appearance, and  $n_H$  is the  
286 harmonic mean population size. The requirement that the mutant phenotype is close to  
287 the resident is the same as above. This result agrees well with that of Ewens (1967),  
288 which gives the probability of fixation as  $2s \frac{n_H}{n(0)}$  in a population with cyclic dynamics.

289 Our result is generalized to a variable, density dependent fitness advantage and arbitrary  
290 ergodic population dynamics. The difference by a factor two is due to different  
291 assumptions on the distribution of surviving offspring – the geometric distribution (used  
292 here), as opposed to the Poisson distribution (as used by Ewens).

### 293 ***The diffusion approximation***

294 The diffusion approximation is very often utilized in population genetics and it can be  
295 used, with care, for the problem of mutant invasion in stochastic environments.

296 Classically, the proportion  $p$  of the invading type is the dynamic state variable and under

297 the assumption that  $p$  changes slowly (between generations) it is sufficient to calculate  
 298 the mean and variance of the change  $\Delta p$  (Kimura 1962). In a stochastic setting, it is  
 299 further necessary to assume that  $p$  changes slowly enough that the full stationary  
 300 distribution of environmental states is experienced during a time-step  $\Delta t$ . Still,  $\Delta t$  has to  
 301 be small enough that  $\Delta p$  is small. In other words, it is required that the invasion process is  
 302 much slower than the stochastic environmental dynamics. Nonetheless, Hill (1972)  
 303 derived the following expressions under the assumptions of discrete generations and  
 304 Poisson distributed offspring:

$$305 \quad \mathbb{E}(\Delta p) = s_A p(1-p) + O(1/n^2) \quad (16a)$$

$$306 \quad \mathbb{V}(\Delta p) = p(1-p)/n_e + O(s_A^2) + O(1/n^2), \quad (16b)$$

307 where  $n_e$  is the harmonic mean population size and  $s_A$  is the arithmetic mean selective  
 308 advantage ( $s_A = \mathbb{E}(e^s - 1) = \bar{s} + O(s^2)$  in our notation). Inserting equations (16a,b) into  
 309 the standard equations of Kimura (1962) yields

$$310 \quad P_0 = (1 - e^{-2n_e s_A / n(0)}) / (1 - e^{-2n_e s_A}), \quad (17)$$

311 expressing the invasion probability of a mutant appearing as a single individual at time 0  
312 ( $p_0 = 1/n(0)$ ). Hill's result has as a first order approximation (assuming  $n_{eSA}$  is large and  
313 discarding terms of order  $s_A^2$  and higher)

$$314 \quad P_0 \approx 2s_A \frac{n_e}{n(0)}, \quad (18)$$

315 which coincides with our result (eq. 15a), apart from the difference in effective  
316 population size.

317 It is possible to derive equations similar to equations (16a,b) also for our models in  
318 discrete and continuous time (not shown). The resulting expressions, similar to equations  
319 (17) and (18), match our results above using the branching process approach (equations  
320 (13a,b) and (15a,b)). In short, it is possible to acquire much the same results using the  
321 diffusion approximation. This is not too surprising, since the necessary assumptions  
322 (large population size, slow invasion) are much the same. However, the conditions under  
323 which the diffusion approximation is valid, especially the averaging across the stationary  
324 distribution of environmental states in equations (16a,b), is somewhat unclear to us. For  
325 example, Hill's (1972) derivation misses the fact that in discrete time, fitness should be  
326 averaged on a logarithmic scale. We leave it to future studies to more thoroughly evaluate  
327 the conditions under which the diffusion approximation is appropriate. Here, we conclude  
328 that it is correct at least to the first order of  $s$ .

329 ***Model examples and tests of accuracy***

330 Continuous time

331 As a continuous time example of our main finding – the probability of mutant invasion –  
332 we choose a theta-logistic model with a birth rate,  $b$ , subject to environmental variation  
333 and a density dependent death rate,  $d$ , according to

334 
$$b(\varepsilon(t)) = d_0 + r + \varepsilon(t) \tag{19a}$$

335 and

336 
$$d(n(t)) = d_0 + r \left( \frac{n(t)}{K} \right)^\theta \tag{19b}$$

337 such that the instantaneous fitness becomes

338 
$$f(n(t), \varepsilon(t)) = b(\varepsilon(t)) - d(n(t)) = r \left( 1 - \left( \frac{n(t)}{K} \right)^\theta \right) + \varepsilon(t). \tag{19c}$$

339  $n(t)$  is the total population size,  $K$  is the carrying capacity, corresponding to the  
340 deterministic equilibrium population size,  $r$  is the per capita growth rate at low densities  
341 and  $\theta$  (together with  $r$ ) controls the shape and strength of density dependence.  $\varepsilon(t)$  is a

342 Gaussian process (more precisely an Ornstein-Uhlenbeck process (Stirzaker 2005)) with  
343 zero mean and an autocovariance function

$$344 \quad \mathbb{C}[\varepsilon(t), \varepsilon(t - \tau)] = \sigma_\varepsilon^2 e^{-|\tau|/T_C}, \quad (20)$$

345 where  $\sigma_\varepsilon^2$  is the stationary variance of the environmental fluctuations and the (auto-)  
346 correlation time  $T_C$  dictates the environmental autocorrelation (the limit  $T_C \rightarrow 0$   
347 corresponds to white noise, with no autocorrelation).

348 As a first example, we choose a resident population with strong density dependence ( $\theta =$   
349  $2$ ) and study the invasion of a mutant with weaker density dependence ( $\theta = 1.98$ ), but the  
350 same equilibrium population size. In the deterministic case ( $\sigma_\varepsilon^2 = 0$ ), the invasion fitness  
351 in this model depends only on the equilibrium population size of the resident,  $K$ ,  
352 compared to that of the invading mutant, and it is a standard result that evolution will  
353 maximize  $K$  (Charlesworth 1971). However, in a variable environment selection will  
354 deviate from the deterministic prediction. The environmental fluctuations have no direct  
355 effect on mean fitness but the resulting fluctuations in population size in combination  
356 with a non-linear density dependence creates selection for weaker density dependence in  
357 this case. This is illustrated in Figure 1, where the density dependent fitness of the  
358 resident ( $f$ , solid, grey line) and the invading mutant ( $\tilde{f}$ , dash-dotted line, mostly  
359 overlapping with  $f$ ) are depicted together with the stationary distribution of resident

360 population size (shaded histogram in background). The fitness difference ( $s = \tilde{f} - f$ , the  
361 thick dashed line is  $100s$ ) is negative for population sizes below  $K$  but positive above  $K$ .  
362 Mean population size is equal to  $K$ , but the strong curvature of  $s$  generates a positive  
363 average fitness advantage for the mutant ( $\bar{s} = 0.0022$ ).

364 We tested the predicted probability of invasion by, first, generating a set of initial  
365 conditions from the stochastic dynamics of the resident population and, next, starting  $10^5$   
366 separate invasion attempts from each initial condition, all initiated from a single mutant  
367 individual (simulation details are given in Appendix B (online supplement)). Figure 2  
368 shows the resulting estimated invasion probabilities plotted against initial population size  
369  $n(0)$  (points with 95% confidence intervals). For the set of parameter values chosen here  
370 (see legend), the results follow our prediction (eq. 13a) very well (dashed line,  $\bar{s}$  and  $n_e$   
371 are calculated from simulations of the population dynamics).

372 We further investigate the robustness of our prediction in Appendix C (online  
373 supplement). To summarize, we find good agreement between our result and more exact  
374 numerical calculations (using eqs. 8a,b) as long as  $n\bar{s}$  is large and  $\bar{s}$  is small. For this  
375 particular model, with these particular parameter values, our approximation has an  
376 average error less than 5% in the region  $50/K < \bar{s} < 0.007$ . At the lower limit,  
377 demographic stochasticity of the resident dynamics is too strong and, more importantly,  
378 the branching process approach is no longer valid since the resident population cannot be  
379 considered infinite from an invasion perspective. Above the higher limit ( $\bar{s} > 0.007$ ), the  
380 variation in  $P_E$  between alternative future environments is too large for our result to hold.

381 In principle, the relationship  $P_0 \sim 1/n_0$  fails. It should here be noted that a diffusion  
382 approximation approach (*sensu* eq. 17) likewise fails at this limit – the difference  
383 between the two predictions is much smaller than the error. We also tested the sensitivity  
384 to strong environmental variation and autocorrelation, and found environmental  
385 autocorrelation to be more critical than variation per se, except close to the boundary  
386 where the risk of extinction of the resident population becomes substantial and the  
387 population undergoes frequent severe bottlenecks. See Appendix C for further details.

388 A technical note: In the derivation of equations (13a,b) we show that for each *possible*  
389 future environment, the probability of invasion converges to the limit as the mean fitness  
390 advantage  $\bar{s}$  goes to zero. Numerical investigations (Appendix C, Figs. C1 and C4) show  
391 that, at least for this model, the *mean* probability, averaged across all possible future  
392 environments, converges much faster than the invasion probabilities corresponding to  
393 single environmental realizations. This means that the value of  $\bar{s}$  may not be as restricted  
394 to really small values as one might conclude from our derivation, and leaves room for  
395 future theoretical investigations on this topic.

#### 396 Discrete time

397 The discrete time example is based on the classical logistic equation, with a fitness (log  
398 per capita growth rate) of the resident population given by

$$399 \quad f(n(t)) = \ln(1 + r(1 - n(t)/K)) + \varepsilon(t) \quad (21)$$

400 We introduce minute amounts of environmental variation here ( $\mathbb{V}[\varepsilon(t)] = 10^{-6}$ ), merely to  
401 avoid completely deterministic dynamics (and loss of ergodicity for some initial  
402 conditions) as we let population size ( $K$ ) grow large.

403 Given stable population dynamics ( $r < 2$ ) and no environmental fluctuations ( $\mathbb{V}[\varepsilon(t)] =$   
404  $0$ ), selection is neutral on the  $r$  parameter. If environmental variation is introduced  
405 through stochastic variation of  $K$  this model generates selection for decreasing  $r$ -values,  
406 basically because a low- $r$  type has weaker density dependence (Turelli & Petry 1980).  
407 The mechanism is very similar to that described in the previous, continuous time example  
408 (Fig. 1). Here we will instead consider the case of unstable dynamics, choosing a high  $r$ -  
409 value, which gives strong, overcompensating density dependence and chaotic dynamics  
410 (in the deterministic case) (May 1974). Selection is still for lower values of  $r$ . To  
411 illustrate several features of our results, we also introduce a trade-off between density  
412 dependence  $r$  and carrying capacity  $K$ , such that a high- $r$  type is compensated with a  
413 higher  $K$ . More precisely, we study the two alternative types 1 and 2:  $\{r_1 = 2.8, K_1 = 10^6\}$   
414 vs.  $\{r_2 = 2.85, K_2 = 1.0023 \cdot 10^6\}$ . Setting type 1 as the resident, type 2 has a fitness  
415 advantage ( $\bar{s} = 0.0023$ ) and can invade (Fig. 3a). On the other hand, if type 2 is  
416 dominating, type 1 has an advantage ( $\bar{s} = 0.0034$ , Fig. 3b). The frequency dependence  
417 comes from the shift in population dynamics as one type or the other dominates the  
418 population. Type 2 has the higher  $r$ -value, which generates more variable population  
419 sizes (compare the distributions of the resident populations in Figs. 3a and 3b). The  
420 strong density fluctuations give type 1 an advantage since it has the lower  $r$ -value.

421 However, once type 1 becomes more common, the population dynamics stabilize  
422 somewhat, such that the advantage is lost. The two types will thus both increase from low  
423 abundances and can coexist in the population. Figure 3c shows a successful invasion of  
424 type 2 (black dots) when type 1 (grey dots) is resident, and the subsequent coexistence.

425 Figures 3a and 3b show a good correspondence between the approximation in equations  
426 (15a,b) and simulation results. In Appendix C we investigate the sensitivity of our  
427 approximation to changes in the resident population size and the strength of selection ( $\bar{s}$   
428 ). We find that the average error is within 5% in the region  $40 / K < \bar{s} < 0.02$ . The upper  
429 limit here is about seven times higher than in the continuous time case, presumably at  
430 least partly due to the fast mixing of the wildly fluctuating dynamics – even rapidly  
431 invading mutants will during the invasion be exposed to a large, representative, portion of  
432 the stationary distribution of the resident type.

433 This example illustrates three things. First of all, that our results are valid for all types of  
434 ergodic dynamics of the resident type (chaos in this case). Secondly, that they are  
435 applicable to situations when invasion does not imply fixation. Thirdly, that population  
436 dynamics may induce frequency dependence. In a constant environment with stable  
437 population dynamics, the feedback environment in the present model is one-dimensional  
438 – it is characterized by a single parameter, the equilibrium population size. If population  
439 sizes fluctuate, on the other hand, the environment in which a new mutant finds itself can  
440 no longer be described so easily – the full distribution of population sizes is necessary to  
441 determine its probability of invasion.

442 *A few technical notes:* This example is not as superficially constructed as it might appear  
443 at first sight. If an  $r$ - $K$  trade-off is modelled as  $r = r_0 + x$  and  $K = K_0(1+cx)$  ( $c > 0$ ), one  
444 quite easily finds parameter values for which there exists an evolutionary branching point  
445 of the trait  $x$  (not shown). In other words, gradual evolution of  $x$  will converge to a  
446 parameter region in which co-existence of closely positioned types is possible (cf. Geritz  
447 et al. 1998). In conclusion, such parameter values are not totally unlikely – they will be  
448 provided by natural selection, given a suitable trade-off. Yet, the model as such is  
449 admittedly superficial and should not be taken too seriously. We choose it here for its  
450 simplicity and the possibility to demonstrate several features of our results with a single  
451 model. Also note that the resident dynamics are strictly speaking not chaotic – the state  
452 space is finite (there can only be a discrete number of individuals) and the dynamics are  
453 stochastic. However, the stochastic dynamics are very similar to the truly chaotic  
454 dynamics of a deterministic, continuous version of the same model.

## 455 **Generalizations**

### 456 Multispecies and multitype evolution

457 The ergodic environment  $E$  can easily be generalized to a community context, or the case  
458 of several coexisting types in a population (or both). As long as the mutant represents a  
459 small phenotypic change of one of the interacting species or one of the coexisting types,  
460 the mean fitness advantage  $\bar{s}$  is well defined and our results are readily applicable. Note

461 that in the multitype case resident population size  $n$  has to be replaced with the number of  
462 individuals of the type from which the mutant descends.

### 463 Diploid, sexual organisms

464 It is likewise straightforward to consider the case of a diploid, randomly mating  
465 population. A new, invading mutant will initially only occur as a heterozygote and its  
466 growth is then equivalent to the asexual case. In continuous time a 'birth event' has to be  
467 interpreted as the event of coupling with a random individual and producing a single  
468 offspring. Each birth event produces a new heterozygote with probability  $1/2$ , which  
469 means the birth rate  $b$  which goes into the equations is the rate of birth events each  
470 heterozygote is involved in divided by two. The assumption of random mating is crucial  
471 here since we cannot allow different mating success for males and females.

472 In the discrete time case the reproductive success of each allele copy needs to follow a  
473 geometric distribution for our analysis to hold. This is for example the case if all  
474 individuals are mated and the number of offspring from each pair of mates has a  
475 geometric distribution (a binomial sampling, due to Mendelian segregation, of a  
476 geometrically distributed number yields a new geometric distribution).

477 In both the discrete and continuous time case it is the mean heterozygote fitness  
478 advantage that enters the equations as  $\bar{s}$ . Completely recessive alleles, which only have  
479 an advantage as a homozygote, are thus not allowed. Further, it is not straightforward to  
480 generalize to the diploid, multitype case, since the multiple genotypes in which a mutant

481 allele may then occur creates an extra source of ‘demographic stochasticity’ not taken  
482 into account here.

### 483 ***Adaptive Dynamics***

484 Given the probability of mutant invasion one can derive expressions for the rate at which  
485 new varieties will invade a population and the consequential rate of trait evolution. If new  
486 types appear as mutants of the resident type with a fixed mutation rate  $\mu$  per individual,  
487 the rate of mutant appearance at any point in time is equal to the number of births times  
488  $\mu$ , which yields the average rate of successful invasions

$$489 \quad \mathbb{E}[\mu b n(0) P_0] \approx 2\mu n_e \bar{s} \quad (22a)$$

490 and

$$491 \quad \mathbb{E}[\mu m(0) P_0] \approx 2\mu n_e \bar{s} \quad (22b)$$

492 in the continuous (eq. 22a) and discrete (eq. 22b) time case, respectively. Note, however,  
493 that  $\mu$  has to be low enough such that only one mutant is invading at any one time. An  
494 immediate conclusion from equations (22a,b) is that evolution is generally slower in  
495 populations with highly variable population sizes, given the same arithmetic mean  
496 population size. This finding is certainly not new, but is here extended to more general  
497 conditions.

498 It is also possible to derive a generalized canonical equation of adaptive dynamics  
 499 (Dieckmann & Law 1996), predicting the rate of evolutionary change over long  
 500 evolutionary time. Considering the evolution of a continuous, heritable trait  $x$  we assume  
 501 the instantaneous fitness of any individual in the population is given by  $f(x_i, E(t))$ , where  
 502  $x_i$  is the trait-value of the individual and  $E(t)$  is the ergodic environment set by a resident  
 503 type with trait value  $x$ . We can then write

$$504 \quad s(t) = f(\tilde{x}, E(t)) - f(x, E(t)) = g(t)\Delta x + O(\Delta x^2)$$

505 where  $\tilde{x}$  is the trait value of a mutant type,

$$506 \quad g(t) = \left. \frac{\partial f}{\partial x_i} \right|_{x_i=x} \quad (23)$$

507 is the instantaneous selection gradient and  $\Delta x = \tilde{x} - x$  is the phenotypic difference in  $x$   
 508 between the mutant and the resident type. Accordingly, we get

$$509 \quad \bar{s} = \bar{g}\Delta x, \quad (24)$$

510 as long as  $\Delta x$  is small, which can be substituted into the expressions for  $P_0$  above.

511 Following much the same procedure as in Dieckmann & Law (1996) gives

512 
$$\frac{dx}{dt} = \mu \sigma_{\mu}^2 n_e \bar{g} \tag{25}$$

513 for both the continuous and discrete time case.  $\mu$  is the mutation rate per individual and  
514  $\sigma_{\mu}^2$  is the variance in mutational effects (on  $x$ ).  $\mu$ ,  $\sigma_{\mu}^2$  and  $n_e$  may all depend on the  
515 resident trait value  $x$ . The effective population size,  $n_e$ , is in the continuous time case  
516 given by equation (13b). In discrete time,  $n_e$  is equal to half the harmonic mean  
517 population size if the assumption of geometrically distributed number of offspring is used  
518 (eq. 15b). A Poisson distributed number of offspring instead yields an effective  
519 population size equal to the harmonic mean population size, using the diffusion  
520 approximation by Hill (eq. 18).

521 Equation (25) seemingly differs by a factor 1/2 from the original expression derived by  
522 Dieckmann & Law for the continuous time case. However, this difference is due to our  
523 definition of effective population size, which converges to  $n/2$  in the deterministic,  
524 continuous time, case. The expression given here has the advantage that it is the same for  
525 discrete and continuous time and that the effective population size in discrete time agrees  
526 with earlier definitions.

527 Equation (25) gives the expected long-term evolutionary change of a continuous trait  $x$ ,  
528 given mutations are of small phenotypic effect and rare, such that consecutive invasions

529 are separated in time. It constitutes a generalized canonical equation of adaptive  
530 dynamics, applicable to arbitrary ergodic environments and population dynamics.

### 531 ***Discussion***

532 We have here calculated the invasion probability of an advantageous mutant type under  
533 quite general conditions. We assume a large, unstructured, monomorphic population and  
534 a mutant of small effect, but put no restrictions on the type of population dynamics or the  
535 variability of the stochastic environment, other than the assumption of ergodicity.  
536 Environmental autocorrelation or slow population dynamics are allowed, as long as  $\bar{s}$  is  
537 small enough that the invasion process is much slower than the population dynamics and  
538 environmental fluctuations. The mutant fitness advantage may depend on population  
539 density as well as environmental conditions. We have outlined how our results can be  
540 generalized to multitype, multispecies scenarios, as well as diploid, sexually reproducing  
541 organisms. We further use the derived invasion probability to calculate the rate of  
542 invasions of new types and to extend the canonical equation of adaptive dynamics, which  
543 shows how our results relate to the rate of adaptation in stochastic environments.

544 The branching process approach used here requires that the average fitness advantage  $\bar{s}$   
545 is small and that  $\bar{sn}$  is large. For theoretical purposes this may not be such a large  
546 problem, but it certainly restricts the number of natural or experimental populations to  
547 which our results can be readily applied. Single invasion experiments in the lab  
548 commonly involve selection coefficients larger than a per cent or two, and experimental

549 populations (except bacteria or protozoans) are commonly too small in numbers. In the  
550 lab or in the field, our predictions can nonetheless serve as benchmark values, in the lack  
551 of a more complete theory. We made some attempts to extend the theory using a  
552 diffusion approximation, but found the results largely conflicted with the same problems  
553 as our first derivation, especially when  $\bar{s}$  is not small. There is still the possibility that  
554 the diffusion approximation does a better job in situations when  $\bar{s}n$  is small to  
555 intermediate – our numerical investigations were not suitable for that type of evaluation –  
556 but a more thorough investigation of the accuracy of the diffusion approximation for this  
557 problem is out of scope here. Moreover, the basis for the application of the diffusion  
558 approximation in this context is in our minds still somewhat shaky and needs further  
559 analysis.

560 Uecker & Hermisson (2011) used an analytical approach very similar to ours (In fact,  
561 their equation (16b) is equivalent to our equation (A4)). However, instead of considering  
562 the stochastic case and taking the limit  $\bar{s} \rightarrow 0$ , Uecker & Hermisson studied a set of  
563 special cases where more complete solutions are attainable – letting the environment or  
564 the resident population change, but in a deterministic fashion. Despite the differences,  
565 many of their conclusions match ours. Among other things, Uecker & Hermisson  
566 demonstrate that in a periodic (sinusoidal) environment, the probability of invasion is  
567 independent of initial conditions if the frequency of environmental change is high enough  
568 (see also Otto & Whitlock 1997). In other words, if the environment changes much faster  
569 than the process of invasion, it is sufficient to take into account the averaged

570 environmental conditions in order to calculate the probability of invasion (save for initial  
571 population size). Further, Uecker & Hermisson demonstrated that the strength of  
572 demographic stochasticity has a direct negative effect on the probability of invasion (our  
573 eq. 13a). It is also possible to show that several of the derived expressions by Uecker &  
574 Hermisson agree with ours if the limit  $\bar{s} \rightarrow 0$  is taken. In our minds, the two studies  
575 complement each other nicely.

576 The importance of the geometric mean fitness, as emphasized in classical bet-hedging  
577 theory, is somewhat implicit in our presentation. In the discrete time case we define  
578 instantaneous fitness  $f$  as the natural logarithm of per capita growth rate, which directly  
579 makes ‘mean fitness’ correspond to the (logarithm of the) geometric mean growth rate.  
580 The classical trade-off between a high arithmetic mean and a low variance is thus not  
581 immediately apparent here, but is incorporated in our definition of ‘fitness’. Instead, the  
582 formalism here emphasizes nonlinearities of the density dependence, sensitivity to  
583 environmental fluctuations and frequency dependence (see also a conceptual discussion  
584 in Ripa *et al.* (2010) on the definition of bet-hedging when fitness is frequency  
585 dependent).

586 In a broader perspective, our results and examples have highlighted several important but  
587 sometimes neglected aspects of trait evolution. Natural populations are subject to  
588 environmental stochasticity, fitness is density and frequency dependent, variable  
589 population sizes induces extra frequency dependence, and the strength or even direction  
590 of selection may differ depending on environmental circumstances. It is our hope that this

591 study may inspire future work towards a more complete theory of trait evolution by  
592 natural selection.

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678

679 **Appendix A**

680 **The limit  $\bar{s} \rightarrow 0$**

681 We here show that

$$682 \quad \lim_{\bar{s} \rightarrow 0} \bar{s} I_E = n(0) \mathbb{E} \left[ \frac{d(E(t))}{n(t)} \right], \quad (\text{A1a})$$

683 where

$$684 \quad I_E = \int_0^\infty \tilde{d}(E(t)) e^{-\int_0^t \tilde{f}(E(\tau)) d\tau} dt \quad (\text{A1b})$$

685 First of all, equation (4) in the main text yields by integration

$$686 \quad n(t) = n(0) e^{\int_0^t f(E(\tau)) d\tau} \quad (\text{A2})$$

687 This implies that

$$688 \quad e^{-\int_0^t \tilde{f}(E(\tau)) d\tau} = e^{-\int_0^t f(E(\tau)) d\tau} e^{-\int_0^t s(E(\tau)) d\tau} = \frac{n(0)}{n(t)} e^{-\int_0^t s(E(\tau)) d\tau}, \quad (\text{A3})$$

689 which substituted into equation (A1b) gives

690 
$$I_E = n(0) \int_0^\infty \frac{\tilde{d}(E(t))}{n(t)} e^{-\int_0^t s(E(\tau)) d\tau} dt = n(0) \int_0^\infty q(t) e^{-\int_0^t s(E(\tau)) d\tau} dt,$$
 (A4)

691 where  $q(t) = \tilde{d}(E(t)) / n(t)$ .

692 The inner integral

693 We need to consider in some detail the behaviour of the inner integral in equation (A4),

694 
$$S(t) = \int_0^t s(E(\tau)) d\tau,$$
 (A5)

695 which is simply a summation of  $s(E(t))$  over time. We assume  $E(t)$  is an ergodic process

696 and we can use the strong or pointwise ergodic theorem (Krengel 1985) to state that for

697 every realization  $E(t)$  and every  $\delta > 0$ , there exists with probability one a  $t_\delta < \infty$  such that

698 
$$1 - \delta < \frac{S(t)}{\bar{s}t} < 1 + \delta, t > t_\delta.$$
 (A6)

699 The relative importance of initial conditions disappear over time, but we note the

700 possibility of realizations  $E(t)$  where the above is not fulfilled, although such possible

701 futures have probability measure zero.

702 We also need to know the behavior of  $t_s$  as  $s$  becomes small, which has to do with how  
 703 fast  $S(t)$  converges to its expectation. It is necessary that  $t_s$  has a finite upper bound in the  
 704 limit  $\bar{s} \rightarrow 0$ . For this we assume, without loss of generality, the mutation corresponds to  
 705 a small change  $\Delta x$  in a heritable trait  $x$  and that the instantaneous fitness advantage has a  
 706 Taylor expansion according to

$$707 \quad s(E(t)) = g(E(t))\Delta x + O(\Delta x^2), \quad (A7)$$

708 where  $g(E(t)) = \frac{\partial s(E(t))}{\partial x}$  is the instantaneous fitness gradient. The limit  $\bar{s} \rightarrow 0$  here  
 709 corresponds to  $\Delta x \rightarrow 0$ . Inserting equation (A7) into equation (A5) gives

$$710 \quad S(t) = \Delta x \int_0^t g(E(\tau)) dt + O(\Delta x^2) = \Delta x G(t) + O(\Delta x^2), \quad (A8)$$

711 where

$$712 \quad G(t) = \int_0^t g(\tau) d\tau. \quad (A9)$$

713 Just like  $S(t)$ ,  $G(t)$  is a simple summation and for every  $\delta > 0$  there exists a finite time  $u_s$

714 such that

715  $(1 - \delta) < \frac{G(t)}{\bar{g}t} < (1 + \delta)$ ,  $t > u_\delta$ . (A10)

716 Note that  $G(t)$ , and thereby  $u_\delta$ , is independent of  $\Delta x$ . Hence,  $u_\delta$  remains fixed (and finite)

717 as we take the limit  $\Delta x \rightarrow 0$  below. Using equations (A8) and (A10) we get

718  $\frac{S(t)}{\bar{s}t} = \frac{\Delta x G(t) + O(\Delta x^2)}{\Delta x \bar{g}t + O(\Delta x^2)} = \frac{G(t)}{\bar{g}t} + O(\Delta x)$ . (A11)

719 and

720  $(1 - \delta) + O(\Delta x) < \frac{S(t)}{\bar{s}t} < (1 + \delta) + O(\Delta x)$ ,  $t > u_\delta$ . (A12)

721 Comparing equations (A6) and (A12) we conclude that for any fixed  $\delta$  we get  $t_\delta \rightarrow u_\delta$  as

722  $\Delta x$  goes to zero.

723 Lower and upper bounds on  $\bar{s}I_E$

724 From equation (A6) it follows that

725  $e^{-\bar{s}t(1+\delta)} < e^{-S(t)} < e^{-\bar{s}t(1-\delta)}$ ,  $t > t_\delta$ , (A13)

726 which can be used to put lower and upper bounds on  $\bar{s}I_E$ :

727  $L_1 + L_2 < \bar{s} I_E < U_1 + U_2,$  (A14a)

728 where

729  $L_1 = n(0)\bar{s} \int_0^{t_\delta} q(t)(e^{-S(t)} - e^{-\bar{s}t(1+\delta)}) dt,$  (A14b)

730  $L_2 = n(0)\bar{s} \int_0^\infty q(t)e^{-\bar{s}t(1+\delta)} dt,$  (A14c)

731  $U_1 = n(0)\bar{s} \int_0^{t_\delta} q(t)(e^{-S(t)} - e^{-\bar{s}t(1-\delta)}) dt,$  (A14d)

732  $U_2 = n(0)\bar{s} \int_0^\infty q(t)e^{-\bar{s}t(1-\delta)} dt.$  (A14e)

733 It is clear that  $L_1$  and  $U_1$  will go to zero as  $\bar{s} \rightarrow 0$ , since we know from above that  $t_\delta$   
 734 remains bounded (it has a finite limit  $u_\delta$  as  $\bar{s} \rightarrow 0$ ).  $L_2$  and  $U_2$  are in principle weighted  
 735 averages of the ergodic process  $q(t)$ , with an exponentially decaying weight function.  
 736 However, as  $\bar{s} \rightarrow 0$  the exponential decay is slower and slower and more and more  
 737 values of  $q(t)$  contribute substantially to the integrals. In short, we use the conjecture that  
 738 integrals of the type

739  $I_c = c \int_0^\infty x(t)e^{-ct} dt$  (A15)

740 go to  $\bar{x}$  as  $c$  goes to zero, as long as  $x(t)$  is ergodic. A formal argument, albeit not a

741 proof, is obtained from the substitution  $\tau = c^{-1}(1 - e^{-c\tau})$ , which gives

$$742 \quad I_c = c \int_0^{1/c} \tilde{x}_c(\tau) d\tau, \quad (A16)$$

743 where  $\tilde{x}_c(\tau) = x(-c^{-1} \ln(1 - c\tau))$  is the process  $x(t)$  with an accelerating time. As  $c$

744 approaches zero the time-transform becomes increasingly linear at lower time-values (a

745 Taylor expansion gives  $-c^{-1} \ln(1 - c\tau) = \tau + \frac{1}{2}c\tau^2 + O(c^2\tau^3)$ ), which supports the

746 conclusion that with probability one

$$747 \quad \lim_{c \rightarrow 0} I_c = \bar{x}. \quad (A17)$$

748 Returning to  $L_2$  (eq. A14c) and  $U_2$  (eq. A14e), we can use equation (A17) to conclude

749 that with probability one

$$750 \quad \lim_{\bar{s} \rightarrow 0} L_2 = \frac{n(0)\bar{q}}{1 + \delta} \quad (A18)$$

751 and

752 
$$\lim_{\bar{s} \rightarrow 0} U_2 = \frac{n(0)\bar{q}}{1 - \delta} . \tag{A19}$$

753 Using equation (A18) and (A19) in equation (A14a) we get

754 
$$n(0)\frac{\bar{q}}{1 + \delta} < \lim_{\bar{s} \rightarrow 0} \bar{s}I_E < n(0)\frac{\bar{q}}{1 - \delta} \tag{A20}$$

755 which is valid for any  $\delta > 0$ . Since we can choose  $\delta$  arbitrarily close to zero we get

756 
$$\lim_{\bar{s} \rightarrow 0} \bar{s}I_E = n(0)\bar{q} \tag{A21}$$

757 with probability one. Returning to the probability of invasion  $P_E$  we have

758 
$$\lim_{\bar{s} \rightarrow 0} \frac{P_E}{\bar{s}} = \lim_{\bar{s} \rightarrow 0} \frac{1}{\bar{s} + \bar{s}I} = \frac{1}{n(0)\bar{q}} \tag{A22}$$

759 for every possible future environment with probability one. Consequently, the

760 expectation of  $P_E$  converges to the same limit, i.e.

761 
$$\lim_{\bar{s} \rightarrow 0} \frac{P_0}{\bar{s}} = \lim_{\bar{s} \rightarrow 0} \frac{\mathbb{E}[P_E | E(0)]}{\bar{s}} = \lim_{\bar{s} \rightarrow 0} \mathbb{E}\left[\frac{P_E}{\bar{s}} | E(0)\right] = \frac{1}{n(0)\bar{q}} \tag{A23}$$

762 Ruling out  $P_E / \bar{s} \rightarrow \infty$

763 Equation (A23) follows from eq. (A22) if we can *completely* rule out the possibility of  
764  $P_E / \bar{s}$  going to infinity. It is thus necessary to show that  $\bar{s}I_E \rightarrow 0$  is not only unlikely,  
765 with probability zero, but *impossible* for all possible future environments  $E(t)$ ,  $t \geq 0$ . For  
766 this, we first write (using the substitution  $T = \bar{s}t$ )

$$767 \quad \bar{s}I_E = n(0)\bar{s} \int_0^\infty q(t)e^{-S(t)} dt = n(0) \int_0^\infty q(T/\bar{s})e^{-S(T/\bar{s})} dT, \quad (\text{A24})$$

768 which in principle behaves as  $n(0) \int_0^\infty q(T/\bar{s})e^{-T} dT$ . It follows that  $\bar{s}I_E \rightarrow 0$  implies the  
769 mutant has, for some unlikely  $E(t)$ , a death rate equal to exactly zero always, or during a  
770 longer-than-zero time-interval an infinite selective advantage  $s$  (such that  $S(t)$  is infinite).  
771 The first options implies a forever immortal mutant, and the second that the mutant has  
772 infinite fitness. We regard both these alternatives as not only unlikely, but impossible (no  
773 organism is immortal and infinite fitness of a small mutation requires a discontinuous  
774 fitness function), which is sufficient for (A23).

775

776 **Figure legends**

777 Figure 1. Instantaneous fitness of the resident type (grey, solid line) and a rare mutant  
778 (dash-dotted dotted line) as functions of the resident population size in the continuous-  
779 time theta-logistic model (eqs. 19a-c), disregarding environmental stochasticity ( $\epsilon$  is set  
780 to 0 when plotting these functions). The dashed line shows the difference between mutant  
781 and resident fitness (x100). The background shading is a histogram (y-scale not shown)  
782 of the population sizes from a simulation of the stochastic resident population dynamics,  
783 where the environmental process is an Ornstein-Uhlenbeck process (Eq. 20) . Parameter  
784 values:  $d_0 = 1$ ,  $r = 1$ ,  $K = 10^6$ ,  $\theta(\text{resident}) = 2$ ,  $\theta(\text{mutant}) = 1.98$ ,  $\sigma_\epsilon^2 = 0.7$ ,  $T_C = 1$ .

785

786 Figure 2. Probability of mutant invasion (y-axis) as a function of the initial resident  
787 population size (x-axis) for the stochastic continuous-time theta-logistic model (eqs. 19a-  
788 c, 20). The black dots (with 95% confidence intervals) indicate the estimated probability  
789 from  $10^5$  simulations, started with a single mutant individual. All invasion attempts for a  
790 given  $n(0)$  were started at the same initial condition. Initial conditions were generated by  
791 simulating the resident population for 100 time units and thereafter until the appropriate  
792 (equally spaced on the log x-axis) population size occurred. The dashed line is the  
793 prediction given by equation (13), where  $\bar{s}$  and  $\overline{(d/n)}$  were calculated from simulations.  
794 The background shading is a histogram of the resident population dynamics, with log-  
795 spaced bins. Parameter values are the same as Figure 1.

796

797 Figure 3. Mutual invasions of two types in the discrete time logistic model (eqs. 6, 7, 21).

798 *a)* Probability of type 2 invading type 1. *b)* Probability of type 1 invading type 2. *a, b)*

799 Estimated invasion probability (black dots with 95% confidence intervals), based on  $10^5$

800 simulations starting at different initial resident population sizes. The dashed line indicates

801 the prediction based on equations (15a,b), where  $\bar{s}$  and  $n_H$  were calculated from

802 simulations. Background shading is a histogram of simulated resident population

803 dynamics (y-scale not shown, but the same in *a)* and *b)*). *c)* A successful invasion of type

804 2 (black dots) into a resident population of type 1 (grey dots). The two types coexisted for

805 at least  $10^4$  generations and showed no signs of one excluding the other (not shown).

806 Parameters, type 1:  $r = 2.8$ ,  $K = 10^6$ ; type 2:  $r = 2.85$ ,  $K = 1.0023 \times 10^6$ .

807

Fig. 1  
Ripa & Dieckmann  
Mutant invasion

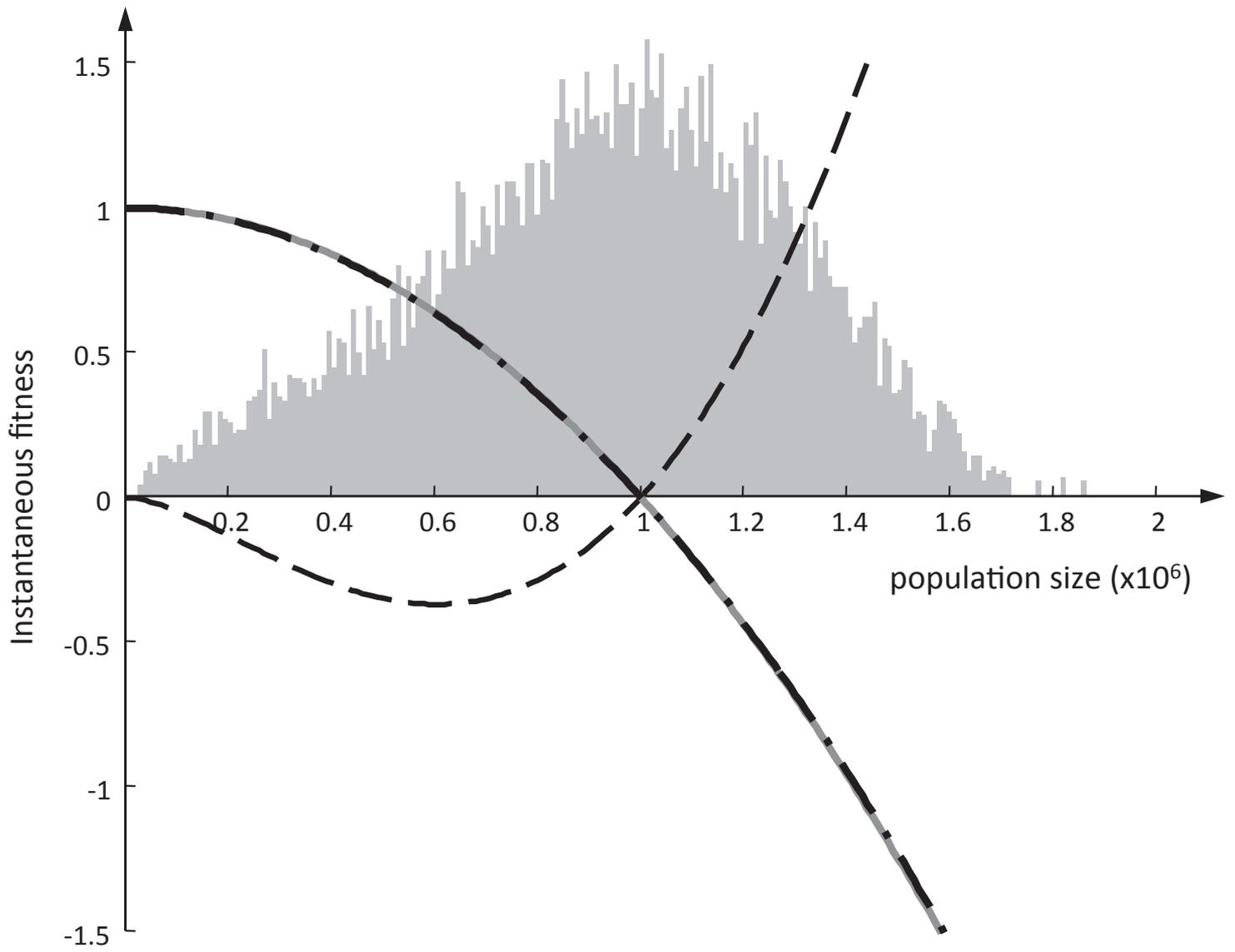


Fig. 2  
Ripa & Dieckmann  
Mutant invasion

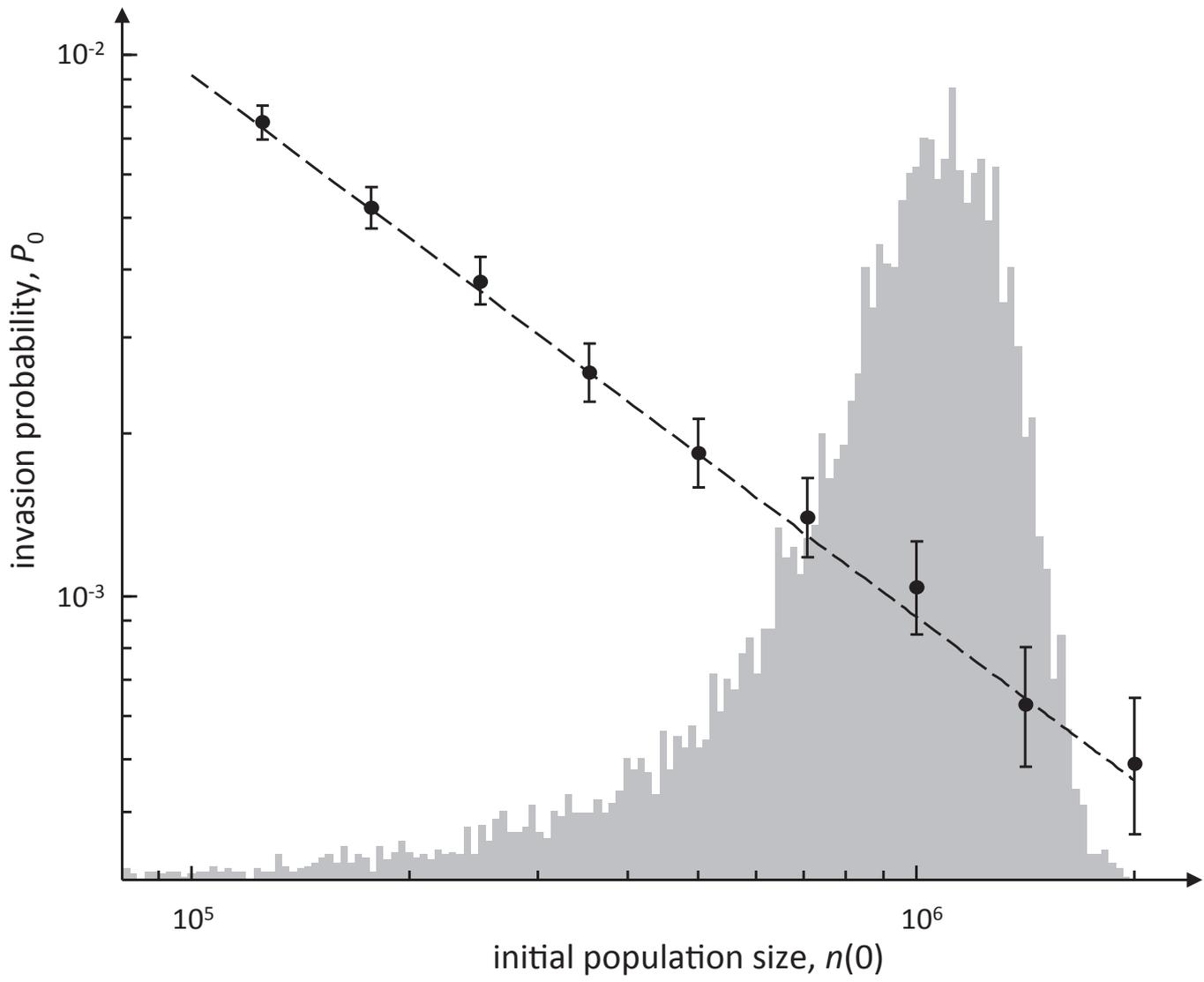
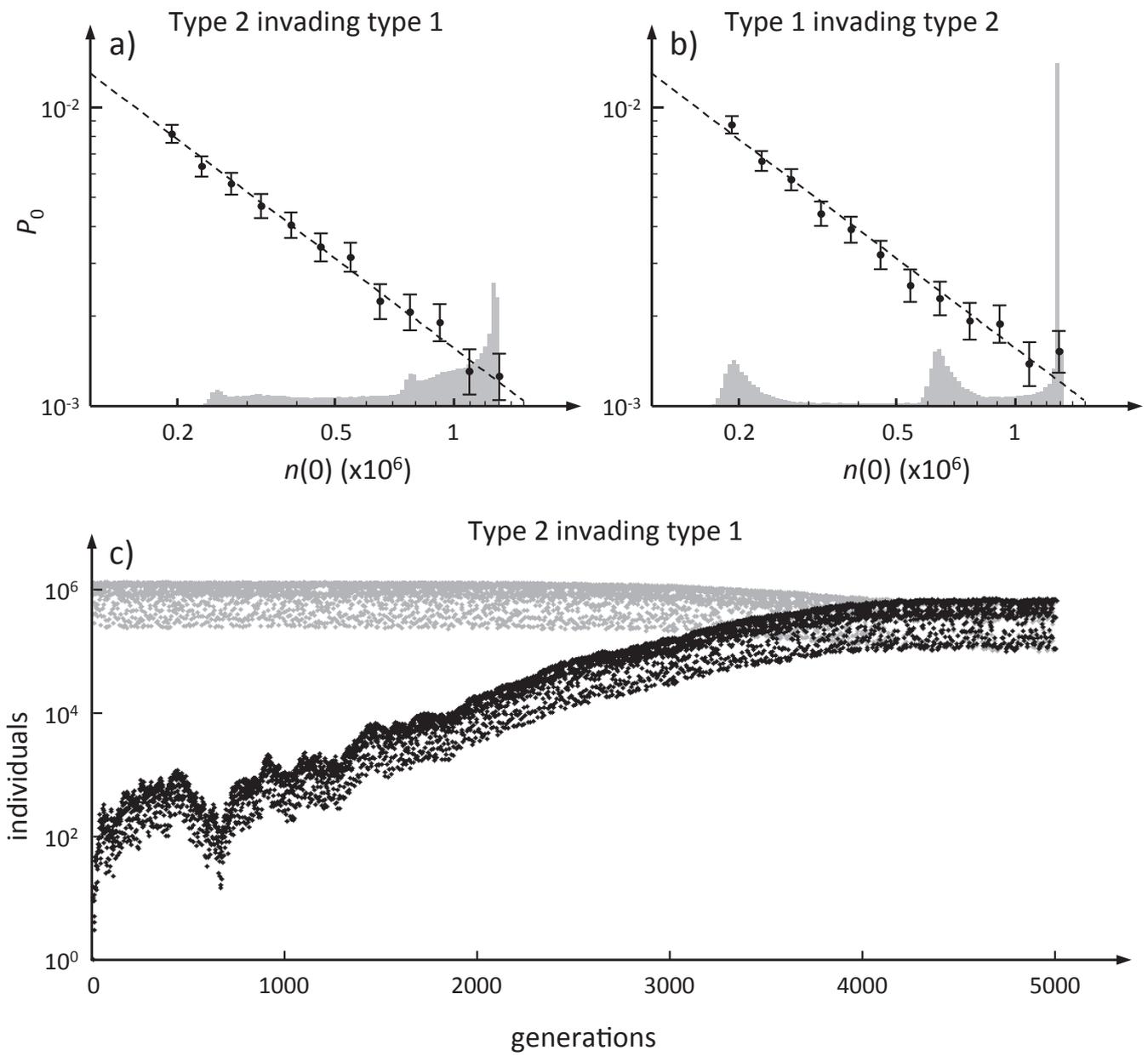


Fig. 3  
Ripa & Dieckmann  
Mutant invasion



## ***Appendix B***

### ***Computer simulation details***

#### Continuous time

The continuous time birth-and-death process was approximated by a discrete time process, with a time interval  $\Delta t$  (a more exact waiting-time approach was in this case too time-consuming). At each time-step, each individual gives birth with probability  $b\Delta t$  and dies with probability  $d\Delta t$ , where the birth and death rates  $b$  and  $d$  depend on the individual's  $\theta$ -value as well as total population size  $n$  and current environmental state  $\varepsilon$  (eqs. 19a,b). Each reproduction produced a new individual identical to the parent.  $\Delta t$  was in the simulations set to  $3.17 \times 10^{-4}$ , chosen such that the total event probability per individual ( $b + d$ ) was equal to 0.001 at equilibrium conditions (Figures 1 and 2). However,  $\Delta t$  was increased to 0.01 in Appendix C to save computer time (This applies to Figures C1, C2 and C3. We also tested  $\Delta t = 0.001$  for a few parameter values, but with no noticeable difference in the results).

The environmental Ornstein-Uhlenbeck process was approximated by a discrete time AR(1) process (Box et al. 1994), with the same autocovariance function (eq. 20). In other words, the environmental process was implemented as

$$\varepsilon_{t+\Delta t} = a\varepsilon_t + v_t, \tag{B1}$$

where

$$a = e^{-\gamma\Delta t} \quad (\text{B2})$$

and  $v_t$  is drawn from a normal distribution with zero mean and variance

$$\mathbb{V}[v] = \sigma_\varepsilon^2(1 - a^2) \quad (\text{B3})$$

Since the discrete time implementation assumes the environment stays constant across a time-step,  $\Delta t$  also has to be small enough that  $\varepsilon_t$  and  $\varepsilon_{t+\Delta t}$  only differ by a small amount, i.e. that the simulation constant  $a$  (eq. B2) is very close to one.

Invasions (Figure 2) were simulated by replacing a single individual of the resident type with an individual of the invading type, and the abundances of the two types were followed over time. A simulation was interrupted as soon as one of them went extinct, and a successful invasion was recorded if the invading type had become fixed.

### Discrete time

At each time-step, each individual was given a geometrically distributed number of offspring (eq. 9), with the mean number of offspring equal to  $e^f$ , where the fitness  $f$  is given by the individual's  $r$  and  $K$  parameters (eq. 21). All parents died after reproduction. A successful invasion was recorded as soon as the invading type had reached an abundance equal to  $K/10$ . At this cutoff point numerical investigations showed that invasion and a long-term coexistence was certain.

All simulations were run in MATLAB® (R2007b, The MathWorks).

## References

Box, G. E. P., G. M. Jenkins and G. C. Reinsel. 1994. Time series analysis: forecasting and control. Prentice-Hall, Upper Saddle River, NJ.

## ***Appendix C***

### ***Tests of accuracy***

We tested the accuracy of the approximate expressions for  $P_0$  in equations (13a,b) (continuous time) and (15a,b) (discrete time) by comparing them to the original expressions for  $P_E$  in equations (8a,b) and (14a,b), averaged across a suite of simulated possible future environments. This was done instead of more explicit simulations of individual invasion attempts, which would be too computer time consuming. We thus rely on the validity of the assumption of an infinite population size, which underlies equations (8a,b) and (14a,b), but gain the ability to investigate larger portions of parameter space.

To estimate the error of equations (13a,b), we used the example continuous time model described in the main text, selected 100 initial conditions from the (simulated) stationary distribution of  $\{n(t), \varepsilon(t)\}$ , and started 50 independent simulations of the resident population dynamics from each initial condition. We then used each simulation to calculate the integrals of equations (8a,b), (Euler method,  $\Delta t = 0.01$ , see Appendix B). We thus acquired 50 measurements of  $P_E$  from each initial condition and calculated their arithmetic mean to get an estimate of  $P_0$ , which was compared to the predicted value

given by equations (13a,b). Figure C1 shows a sample of estimated  $P_E$ -values (black dots) together with the corresponding estimated  $P_0$ -values (red crosses), our prediction (eqs. 13a,b, blue lines) and a diffusion approximation (eq. 17, green lines).

The error in the predicted  $P_0$  for each initial condition  $i$  was calculated as  $e_i = \log((\text{predicted } P_0)/(\text{estimated } P_0))$ , and the total error for each parameter setting was measured as the square root of the bias-corrected mean squared error, according to

$$e_{tot} = \sqrt{\frac{1}{100} \sum_{i=1}^{100} e_i^2 - cv^2 \mu_0 - cv^2}, \quad (\text{C1})$$

where  $cv^2 = \frac{1}{100} \sum_{i=1}^{100} \frac{s_{P,i}^2}{\hat{P}_i^2}$  is the mean squared relative standard error,  $s_{P,i}^2$  is the squared standard error of the estimated  $P_0$  for initial condition  $i$ , and  $\mu_0$  is the (estimated) mean prediction error, across initial conditions. The bias correction is based on the assumption that  $P_E$  has a constant coefficient of variation, independent of initial conditions, and Taylor expansions of the log transform. Qualitatively, our results are the same, with or without the bias correction. The error estimate in equation (C1) can be interpreted as the mean relative error of our prediction, averaged across initial conditions. It includes a possible constant bias ( $\mu_0$ ) as well as variation between initial conditions not captured by the predicted  $1/n(0)$  relationship (eq. 13a).

Figure C2 shows the estimated relative error (eq. C1) for different values of the population carrying capacity,  $K$ , and the mean fitness advantage,  $\bar{s}$ . The calculations are, to be precise, carried out for constant values of  $\Delta\theta$  ( $-2.5 < \log_{10}(-\Delta\theta) < 0$ ), and the

corresponding  $\bar{s}$  varies somewhat depending on the value of  $K$ . This variation is, however, very small and a correction for this would not change any conclusions drawn from Figure C2. The greyscale shading and solid line contour levels depict the estimated error. The dashed straight lines indicate the boundaries  $50/K < \bar{s} < 0.007$ , which approximates the region where the error is less than 5%.

Figure C3 shows the dependence of the error on the variance ( $\sigma_\epsilon^2$ ,  $x$ -axis) and correlation time ( $T_C$ ,  $y$ -axis) of the environmental fluctuations (see eq. 20). The invading mutant has a  $\theta$ -value of 1.995 (compared to the resident  $\theta = 2$ ), but the different environmental parameters would generate different values of  $\bar{s}$ , all else being equal. For a fair comparison between different values of  $\sigma_\epsilon^2$  and  $T_C$ , we adjusted the mutant  $K$ -value such that the mutant has a fixed average fitness advantage  $\bar{s} = 0.0002$ . This  $K$ -adjustment is always small (less than  $10^{-5}K$ ) and shifts sign from positive at low values of  $\sigma_\epsilon^2$  to negative at high values of  $\sigma_\epsilon^2$ . Further, the initial conditions are always the same 50 conditions sampled from the stationary distribution of the standard parameter values  $\sigma_\epsilon^2 = 0.4$ ,  $T_C = 1$ .

The error depicted in Figure C3 is large at high values of  $T_C$  and close to the region where the resident population goes extinct too quickly for measurements to be possible (dotted region). That our approximation fails in slowly fluctuating environments (a large  $T_C$ ) is not surprising, since one of the main assumptions is that the environmental fluctuations are much faster than the invasion process. This is confirmed by trial calculations with ten times faster invasions ( $\bar{s} = 0.002$ ), which basically shifts the error contour levels to ten

times lower values of  $T_C$  (not shown). When the population dynamics are very violent, close to the dotted region in Figure C3, a close inspection of the population dynamics shows that the resident population goes through repeated periods of very low densities, several orders of magnitude below  $K$ . Each such bottleneck of the resident population strikes the mutant too, since they are ecologically very similar, and has a large negative impact on the probability of invasion. The total probability becomes highly dependent on the exact number of bottlenecks during an invasion, which causes a large variation in invasion probability between different realizations of the environmental process, despite a very long invasion time. It follows that the assumptions of our derivation are not fulfilled and the approximation fails (it requires an even smaller value of  $\bar{s}$ ).

Figures C4-5 show the same calculations as Figures C1-2, but for the discrete time model (eq. 21). In figure C5 it can be seen that the region where the error is less than 5% is now larger ( $40/K < \bar{s} < 0.019$ ), especially at the upper end. The reason for this is hard to disentangle completely, but one answer might be the rapid chaotic fluctuations of population size in this model, which means an invading mutant is quickly exposed to the full range of environmental fluctuations. This model is also, at least in the short term, much more deterministic than the continuous time model. The resident population sizes during the important first few generations after the first appearance of a new mutant are highly predictable, given the initial population size. There is thus relatively little variation between different realizations of  $I_E$  (there is a relatively small spread of black dots in Figure C4), which reduces the possible error related to taking the mean of a function as the function of the mean ( $P_0$  is the mean of  $P_E$ , which is a non-linear function of  $I_E$  (eq. 14a)). Finally, we would like to point out that a diffusion approximation succeeds within

almost exactly the same region of parameter space, a region within which the difference between the two predictions is still small.

## Figure Legends

Figure C1. Samples of simulated invasion probabilities,  $P_E$  (y-axis, eq. 8a), of the continuous time example model (eqs. 19, 20). For each parameter setting (panel), 100 initial conditions  $\{n(0), \varepsilon(0)\}$  were chosen from the simulated stationary distribution of  $\{n(t), \varepsilon(t)\}$  and for each initial condition the future population dynamics was simulated 50 times to give 50 estimates of the conditional invasion probability  $P_E$  (black dots, eq. 8a). Red crosses: The estimated unconditioned invasion probability,  $P_0$ , calculated as the arithmetic mean of the  $P_E$ -values (eq. 9). Blue lines: The predicted  $P_0$  (eq. 13a). Green dashed lines (often coinciding with the blue lines): The diffusion approximation (eq. 17 with  $n_e$  from eq. 13b). Black dash-dotted lines: The neutral prediction ( $1/n(0)$ ). Resident population parameters:  $d_0 = 1$ ,  $r = 1$ ,  $\theta = 2$ ,  $\mathbb{V}(\varepsilon_t) = \sigma_\varepsilon^2 = 0.4$ ,  $T_C = 1$ . The carrying capacity  $K$  differs between the panel rows and is indicated in the left hand margin. The invading type has a  $\theta$ -value equal to  $2 - \Delta\theta$ , where  $\Delta\theta = 0.00316, 0.0178, 0.178$  and  $1.00$  in the panel columns, left to right, respectively. The corresponding mean fitness advantage,  $\bar{s}$ , is indicated on the top of each column (the dependence on  $K$  is small, less than 2%).

Figure C2. The average relative error (eq. C1) of the predicted  $P_0$  (eq. 13a) (grey shading and contour lines), depicted as a function of the mean fitness advantage  $\bar{s}$  ( $x$ -axis) and the carrying capacity  $K$  ( $y$ -axis). The effective population size  $n_e$  (eq. 13b) is approximately  $0.24K$ . The region  $50/K < \bar{s} < 0.007$ , roughly where the error is less than 5%, is indicated by a black dashed line. The figure is based on a grid of 11  $\Delta\theta$ -values and

12  $K$ -values, equally spaced on a logarithmic scale (see Fig. C1 and the main text for further details).

Figure C3. The average relative error (eq. C1) of the predicted  $P_0$  (eq. 13a) (grey shading and black contour lines), depicted as a function of the variance ( $x$ -axis) and correlation time ( $y$ -axis) of the external environment ( $\varepsilon_t$ , see eqs. 19, 20). The model and most parameter values are as in Fig C1. The resident has  $K = 10^8$  and  $\theta = 2$ . The invading type has  $\theta = 1.995$  and a  $K$ -value adjusted such that  $\bar{s} = 0.0002$ , irrespective of strength and autocorrelation of the environmental fluctuations. In the dotted area, the extinction rate of the resident population was too high for meaningful measurements.

Figure C4. Same as Figure C1, but for the discrete time model (eq. 21). The resident type has  $r = 2.8$  (corresponding to chaotic dynamics) and the invading mutant has  $r = 2.8 - \Delta r$ , where  $\Delta r$  ranges from  $10^{-2.5}$  to 1, equally spaced on a logarithmic scale, in steps of  $10^{0.5}$ . The  $K$ -values are spaced similarly, from  $10^4$  to  $10^8$ . At  $K$ -values below  $10^4$ , the resident population went extinct too quickly. Only a sample of the simulation results are depicted here. Black dots:  $P_E$ -values (eq. 14a). Red crosses:  $P_0$  (mean  $P_E$ ). Blue lines: predicted  $P_0$  (eqs. 15a,b). Green dashed lines: diffusion approximation (eq. 17 with  $n_e$  given by eq. 15b). Black dash-dotted lines: The neutral prediction ( $1/n(0)$ ). Each row of panels corresponds to a fixed value of  $K$ , as indicated in the left margin. Each column corresponds to  $\Delta r = 0.00316, 0.0316, \text{ and } 1.00$ , from left to right, respectively. The

corresponding mean fitness advantage,  $\bar{s}$ , is indicated on the top of each column (the dependence on  $K$  is small, less than 2%).

Figure C5. The average relative error (eq. C1) of the predicted  $P_0$  of the discrete time model (eq. 21), depicted as a function of the mean fitness advantage  $\bar{s}$  ( $x$ -axis) and the carrying capacity  $K$  ( $y$ -axis). The effective population size  $n_e$  (half the harmonic mean) is roughly  $0.34K$ . Other details are given in Figure C4 and Appendix C. The dashed lines mark the boundaries of the region  $40 / K < \bar{s} < 0.019$ , where the mean relative error is below 5%.

Fig. C1  
Ripa & Dieckmann  
Mutant invasion

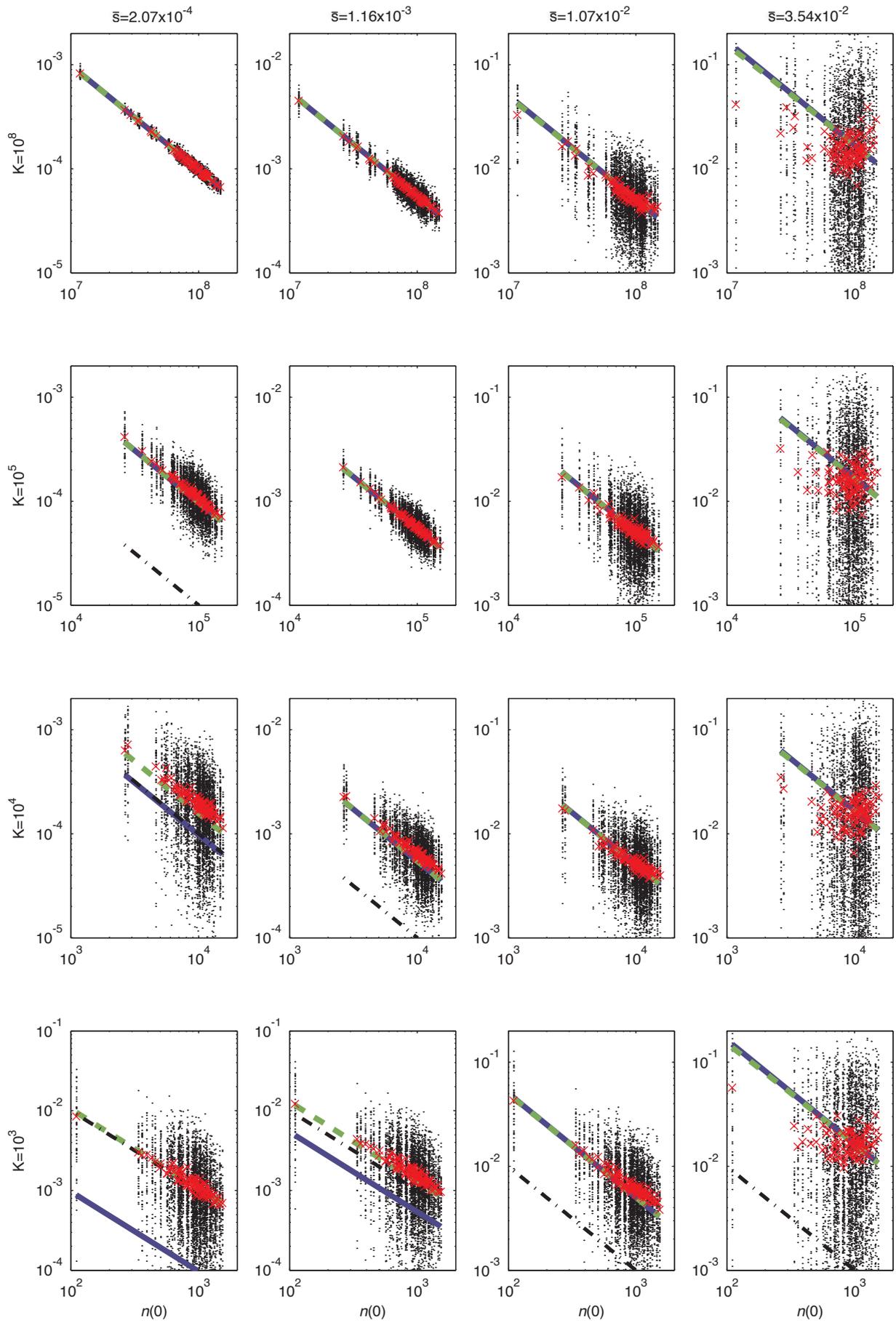


Fig. C2  
Ripa & Dieckmann  
Mutant invasion

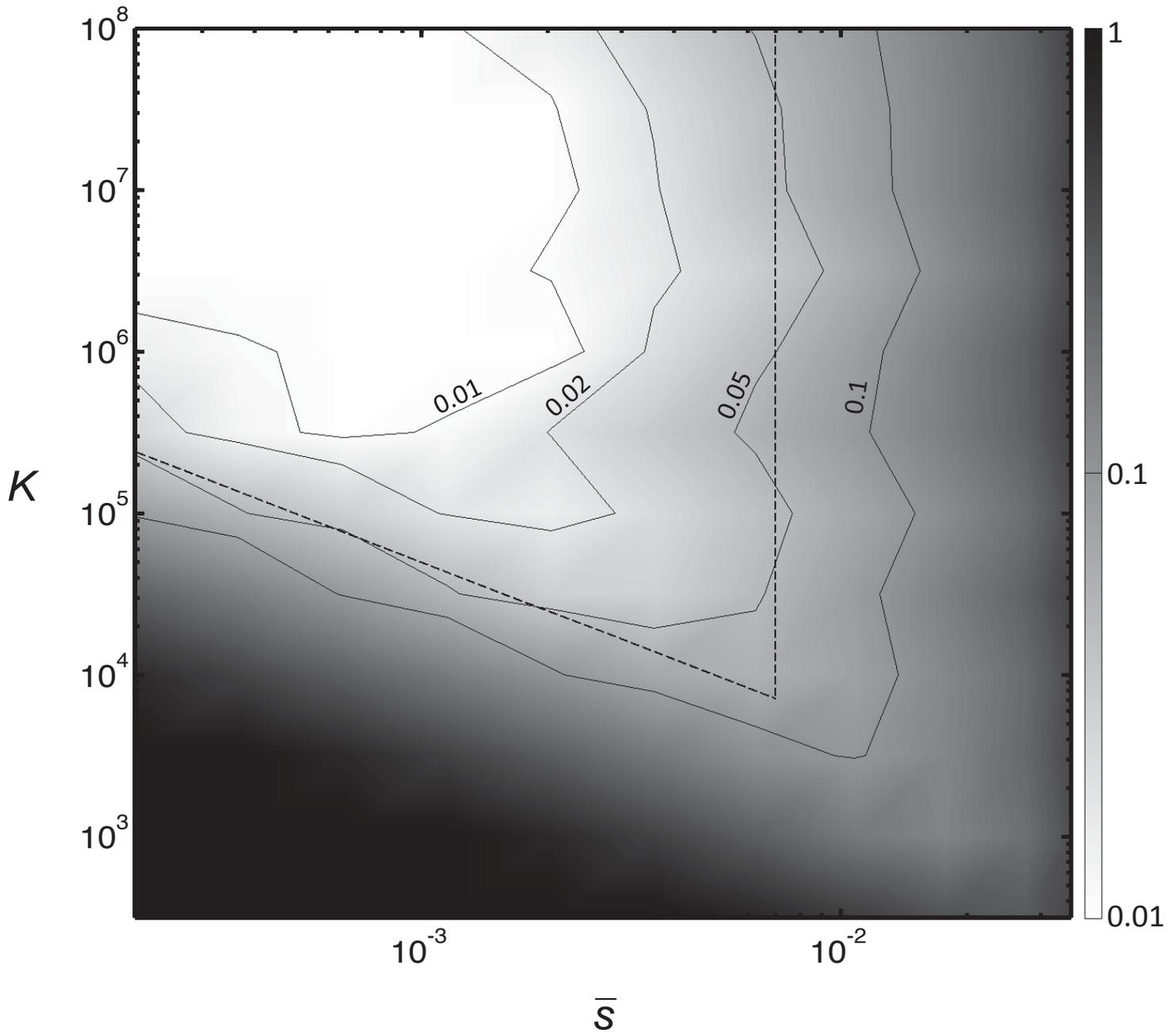


Fig. C3  
Ripa & Dieckmann  
Mutant invasion

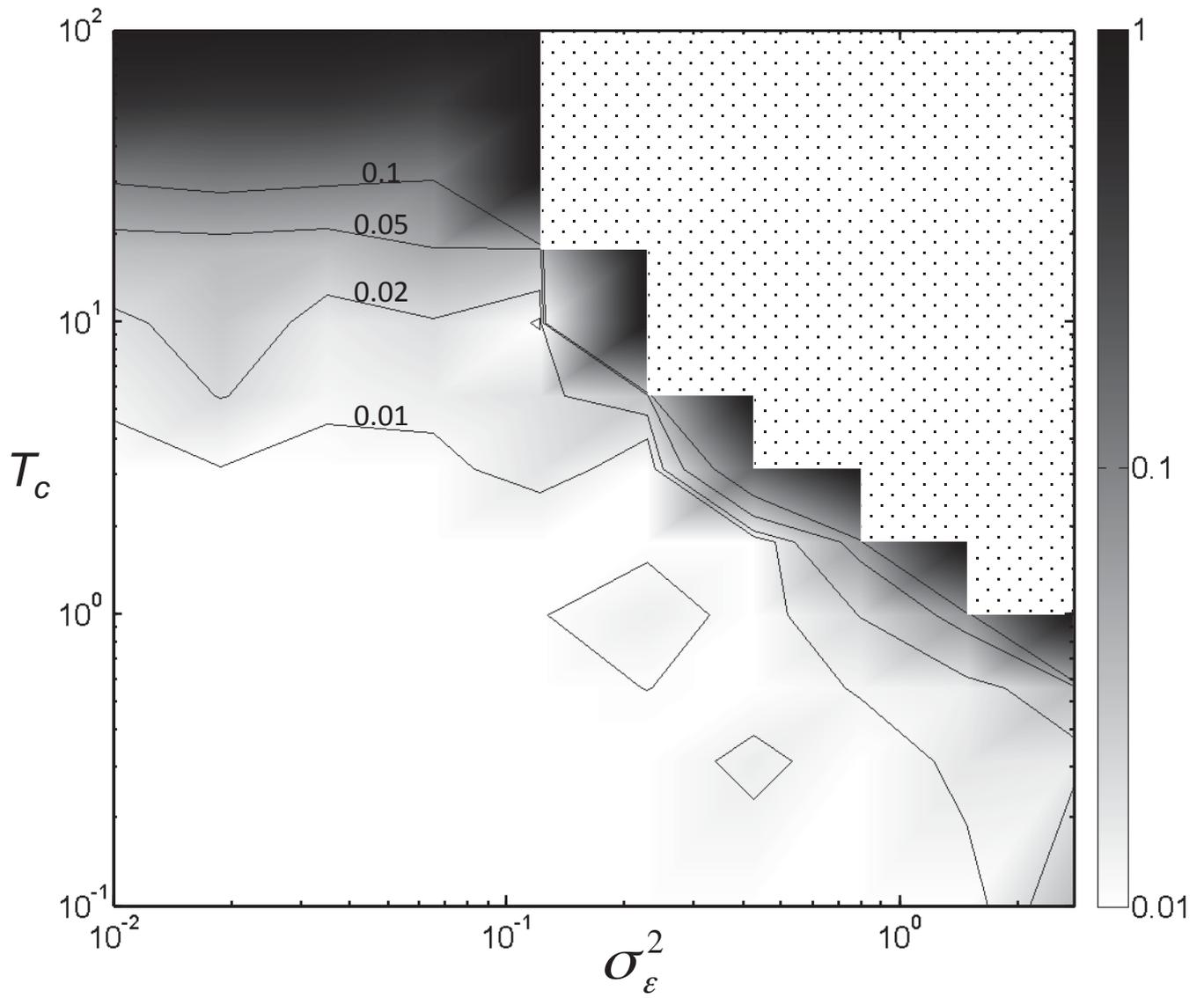


Fig. C4  
Ripa & Dieckmann  
Mutant invasion

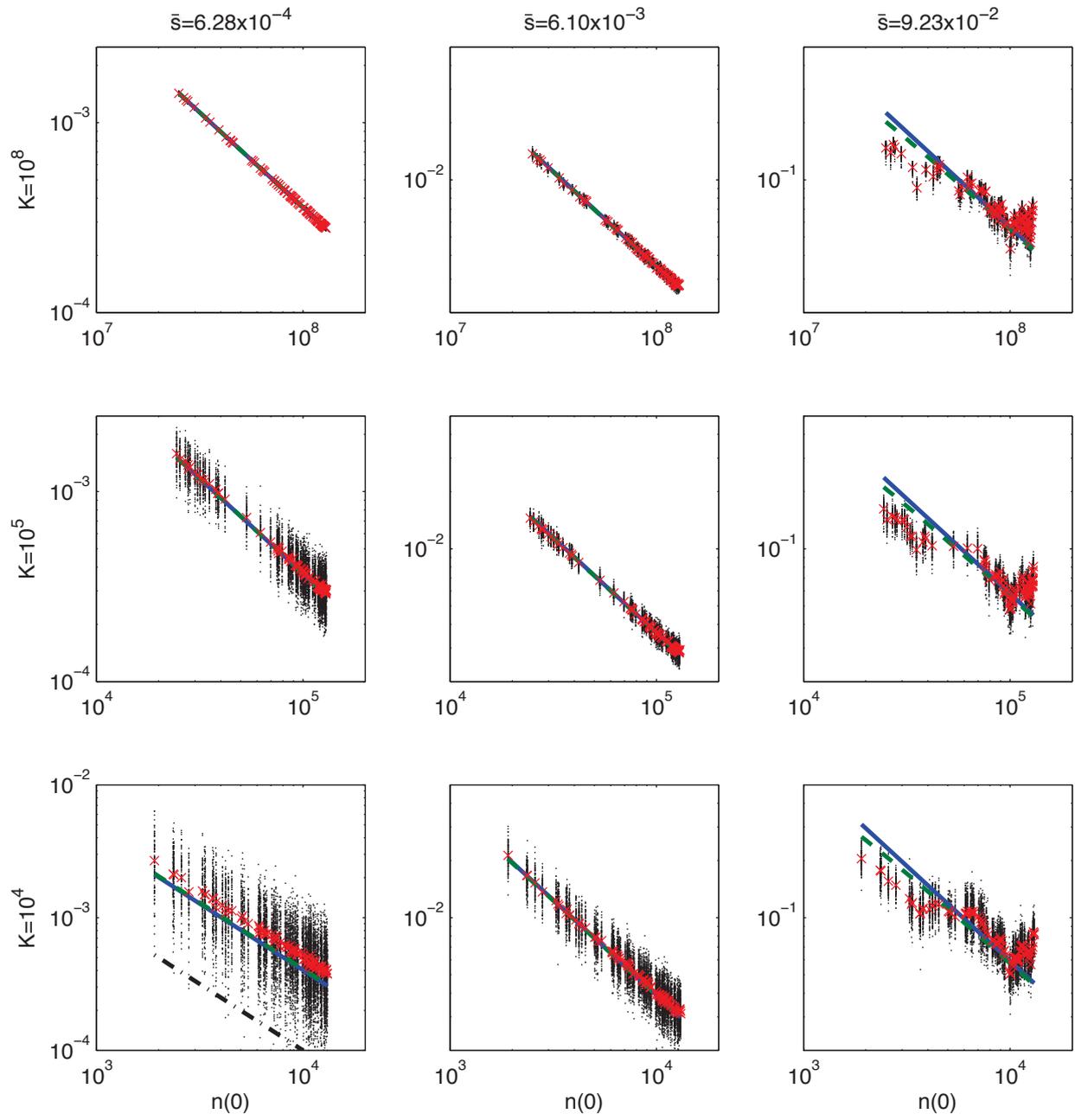


Fig. C5  
Ripa & Dieckmann  
Mutant invasion

