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### A derivation of the statistical characteristics of forest fires

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# A DERIVATION OF THE STATISTICAL CHARACTERISTICS OF FOREST FIRES

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## Abstract

2

The analysis of large data sets concerning fires in various  
3 forested areas of the world has pointed out that burned areas  
4 can often be described by different power-law distributions for  
5 small, medium and large fires and that a scaling law for the time  
6 intervals separating successive fires is fulfilled. The attempts of  
7 deriving such statistical laws from purely theoretical arguments  
8 have not been fully successful so far, most likely because im-  
9 portant physical and/or biological factors controlling forest fires  
10 were not taken into account. By contrast, the two-layer spatially  
11 extended forest model we propose in this paper encapsulates the  
12 main characteristics of vegetational growth and fire ignition and  
13 propagation, and supports the empirically discovered statistical  
14 laws. Since the model is fully deterministic and spatially ho-  
15 mogeneous, the emergence of the power and scaling laws does  
16 not seem to necessarily require meteorological randomness and  
17 geophysical heterogeneity, although these factors certainly am-  
18 plify the chaoticity of the fires. Moreover, the analysis suggests  
19 that the existence of different power-laws for fires of various scale  
20 might be due to the two-layer structure of the forest which allows  
21 the formation of different kinds of fires, i.e. surface, crown, and  
22 mixed fires.

23

24      *Keywords:* Forest fires; Wildfire statistics; Model; Vegetational  
 25      growth; Power law; Scaling law

26      **Empirical evidence of forest fires characteristics**

27      Forest fires have been observed for centuries all over the world, and huge  
 28      data sets are now of public domain. They usually contain long series of  
 29      fire events identified by location, time of occurrence, and burned area.  
 30      Statistical analyses of these data sets have allowed various authors to  
 31      identify, on a purely empirical basis, general characteristics of forest  
 32      fires.

33      Malamud et al. (1998) and Ricotta et al. (1999) were the first to  
 34      perform statistics of the burned areas. They arrived to the same con-  
 35      clusion, namely that burned areas are distributed as a power law, rep-  
 36      resented by a straight line in log-log scale. This conclusion is actually  
 37      surprising, because the only graph reported in Ricotta et al. (1999)  
 38      clearly shows that the distributions of small, medium and large fires  
 39      are well approximated by different power laws, and the same, though  
 40      less pronounced, effect is detectable in the plots obtained by Malamud  
 41      et al. (1998). Most likely, this slightly distorted interpretation of the  
 42      results had two targets: find an agreement with the theoretical studies  
 43      available at that time on self-organized critical forest-fire models (see  
 44      next section), and support the idea that the knowledge of the occur-  
 45      rence frequency of small and medium fires can be used to quantify the  
 46      risk of large fires.

47                        [Figure 1 about here.]

48      Subsequent studies (Ricotta et al. (2001); Song et al. (2001); Reed  
 49      and McKelvey (2002) and, in particular, Ricotta (2003)) confirmed that  
 50      the distributions of the burned areas are smooth but can sometimes be  
 51      approximated by three or two different power laws, as shown in Fig. 1,  
 52      where three examples taken from the literature are reported.

53      A few years later, the first studies on the temporal distributions  
 54      of the fire events are performed through various statistical techniques  
 55      (Telesca et al. (2005); Lasaponara et al. (2005); Ricotta et al. (2006)).

56 The main result is the discovery of a high degree of time-clusterization  
 57 even if the burned areas are not distributed as a power-law. This means  
 58 that the occurrence of large events mimics the process of occurrence of  
 59 smaller events, thus allowing one to model the scarce big fires on the  
 60 basis of the abundant small fires. This is neatly pointed out in Corral  
 61 et al. (2008) where the fire catalog for all Italy in the period 1998-  
 62 2002 is used to estimate the probability density  $D(\tau|s)$  of the time  
 63 intervals  $\tau$  separating two successive fires within the so-called class-  
 64  $s$  fires (i.e., fires with burned areas grater than or equal to  $s$ ). The  
 65 distributions estimated for each class (see the curves displayed in Fig.  
 66 2(a) reproduced from Corral et al. (2008)) can somehow be fitted with  
 67 a power-law, but the exponent of the power-law (i.e., the negative slope  
 68 of the curve) decreases with the increase of the minimum burned area  
 69  $s$  characterizing the class. However, all these distributions practically  
 70 collapse into a single function  $F$ , as shown in Fig. 2(b) (again extracted  
 71 from Corral et al. (2008)), through the simple scale transformation  
 72  $\tau \rightarrow R(s)\tau$  and  $D(\tau|s) \rightarrow D(\tau|s)/R(s)$ , where  $R(s)$  is the rate of fire  
 73 occurrence in class  $s$  (defined as the mean number of fires per unit time  
 74 with burned area greater than or equal to  $s$ ). This interesting discovery,  
 75 formally revealed by the relationship

$$D(\tau|s) = R(s)F(R(s)\tau) \quad (1)$$

76 allows one to conclude that forest fires fulfill a scaling law for the time in-  
 77 tervals separating successive fires without necessarily displaying power-  
 78 law distributions of the burned areas.

79 [Figure 2 about here.]

#### 80 **Theoretical investigations on forest fire characteristics**

81 The attempts of deriving fire characteristics from purely theoretical  
 82 arguments have been performed through two different classes of models.  
 83 The models of the first class, known as self-organized critical forest-fire  
 84 (SOCFF) models, are probabilistic cellular automata defined over a  
 85 square lattice with  $L^2$  sites. In the first version (Bak et al. (1990))  
 86 each site is at each time step in one of three possible states: green

87 (i.e., not burning) tree; red (i.e., burning) tree; absence of vegetation.  
88 The transition rules are very simple: (*i*) green trees become red if they  
89 are close to red trees and remain green otherwise; (*ii*) red trees die  
90 thus leaving the site empty; (*iii*) each empty site has a probability  $p$  of  
91 becoming occupied by a green tree. Bak et al. (1990) state that their  
92 model is a self-organized critical model capable of showing how the solar  
93 energy absorbed continuously at low rate by vegetation can be randomly  
94 dissipated through rare and disruptive events (the fires). However, the  
95 agreement with real forests is, even qualitatively, rather poor because  
96 the model generated fires are always present (in the form of travelling  
97 fronts burning pieces of the boundaries of vegetational clusters).

98 The model proposed by Bak et al. (1990) is immediately criticized  
99 by Drossel and Schwabl (1992) who point out some of its critical aspects  
100 and introduce a second parameter ( $f$ ), called "lightning parameter" by  
101 means of which they modify rule (*i*) saying that green trees not close to  
102 red ones have a probability  $f$  to become red. This variation introduces  
103 a random exogenous mechanism of fire ignition and is essential for cre-  
104 ating clusters of fires with areas distributed as power-laws. A number  
105 of variants of the SOCFF model are immediately proposed by various  
106 authors (see Clar et al. (1996) for a review). In particular, Drossel and  
107 Schwabl (1993) introduce a third parameter, called "tree immunity" in  
108 order to modify, once more, rule (*i*) by saying that green trees have a  
109 certain probability of remaining such when they are close to red trees.  
110 Later (Song et al. (2001)) this variant is shown to give rise to distribu-  
111 tions of burned areas that can be approximated with two power-laws,  
112 one for small-medium fires and one for large fires. A similar result is  
113 obtained by Schenk et al. (2000) by stressing the finite-size effects in  
114 SOCFF models.

115 In the second class of models, here called two-layer models, the  
116 forest is described by two sets of ordinary differential equations, one  
117 associated with the lower layer composed of bryophytes, herbs, shrubs  
118 or any mix of these plants and the other associated with the upper layer  
119 composed of plants and trees of various species. The growth of the two  
120 layers in the absence of fire is described in the standard continuous time

121 form

$$\begin{aligned}\dot{L} &= r_L L \left(1 - \frac{L}{K_L}\right) - \alpha L U \\ \dot{U} &= r_U U \left(1 - \frac{U}{K_U}\right)\end{aligned}\tag{2}$$

122 where  $r$  and  $K$  indicate growth rate and carrying capacity and  $\alpha L U$   
123 is the surplus of mortality in the lower layer due to light interception  
124 caused by tree canopy. Thus, in the absence of fire, trees grow logisti-  
125 cally toward the carrying capacity  $K_U$ , while plants of the lower layer  
126 tend toward  $(1 - \alpha K_U / r_L) K_L$ . The validity and limitations of eq. (2)  
127 are discussed in Casagrandi and Rinaldi (1999), where realistic values of  
128 the five vegetational parameters ( $r_L, r_U, K_L, K_U, \alpha$ ) are also suggested.

129 As for the fire, there are two options. The first (Casagrandi and  
130 Rinaldi (1999)) is to add two extra-variables representing the burning  
131 (red) biomasses in the two layers and describe the propagation of the  
132 fire to the green biomasses  $L$  and  $U$  through suitable fire attack rates.  
133 This gives a model with four ordinary differential equations which is,  
134 however, a so-called slow-fast model because the green biomasses grow  
135 very slowly (typically over years), while the red ones become suddenly  
136 very high when the fire starts and then practically drop to zero after a  
137 very short time (typically a few days or weeks).

138 The second option (Maggi and Rinaldi (2006)) is to push the slow-  
139 fast nature of the system to the extreme, by considering fires as devas-  
140 tating events capable of reducing instantaneously the green biomasses  
141 of finite amounts. This can be accomplished, without adding extra  
142 differential equations, by defining as shown in Fig. 3(a) the pre- and  
143 post-fire manifolds  $\mathcal{X}^-$  and  $\mathcal{X}^+$  and the map from  $\mathcal{X}^-$  to  $\mathcal{X}^+$  interpreting  
144 the impact of the fire.

145 [Figure 3 about here.]

146 The pre-fire manifold  $\mathcal{X}^-$  in Fig. 3(a) is piece-wise linear and non-  
147 increasing, and the set below the manifold is convex. The first property  
148 is obvious because less fuel originated from trees (i.e. less trees) is nec-  
149 essary for fire ignition if more fuel originated from bushes is available  
150 on the ground. The second property simply says that if  $x' = (B', T')$   
151 and  $x'' = (B'', T'')$  are two states of the forest at which fire ignition is

not possible (i.e. two points below the manifold  $\mathcal{X}^-$ ) no mix of these two states (i.e. no points of the segment connecting  $x'$  with  $x''$ ) can give rise to fire ignition. A formal support of these two properties can be found in Maggi and Rinaldi (2006). The geometry of the pre-fire manifold allows one to sharply identify surface fires (vertical segment of  $\mathcal{X}^-$ ), crown fires (horizontal segment of  $\mathcal{X}^-$ ) and mixed fires (oblique segment of  $\mathcal{X}^-$ ). By definition, surface fires do not involve the upper layer, so that the post-fire conditions are on the vertical segment characterized by  $L^+ = \lambda_L \rho_L K_L = \lambda_L L^-$ . In other words,  $\rho_L$  is, by definition, the portion of the lower layer carrying capacity  $K_L$  at which surface fires occur and  $\lambda_L$  is the portion of the lower layer biomass that survives to surface fire. Similarly, fires in the upper layer are characterized by a vertical instantaneous transition from  $U^- = \rho_U K_U$  to  $U^+ = \lambda_U U^-$ . The most extreme surface fire is represented by the transition  $S^- \rightarrow S^+$ , while the most extreme crown fire is represented by the transition  $C^- \rightarrow C^+$ . The assumption that mixed fires initiate on the segment  $C^- S^-$  implies, by continuity, that post-fire conditions are on a curve connecting  $C^+$  and  $S^+$  which, for simplicity, is identified with the linear segment  $C^+ S^+$ .

Fire sequences can be easily obtained from the model, as shown in Fig. 3(b). Starting from a given initial condition, say point 0, one numerically integrates the differential eqs. (2) until the solution hits the pre-fire manifold  $\mathcal{X}^-$  at point  $1^-$ . Then, using the map  $\mathcal{X}^- \rightarrow \mathcal{X}^+$  one can determine the post-fire conditions, namely point  $1^+$ . Finally, the procedure is iterated and a series of fires  $(2^- \rightarrow 2^+), (3^- \rightarrow 3^+), \dots$  is obtained.

A detailed analysis of this minimal model (Maggi and Rinaldi (2006)) has shown that it is very flexible and can reproduce, by tuning its parameters, the fire regimes of savannas, boreal forests and Mediterranean forests. The dependence of the model behavior upon its numerous parameters has been thoroughly investigated in Dercole and Maggi (2005) and in Bizzarri et al. (2008). Moreover, long series of model generated fires have been statistically analyzed and the result is that the distributions of the total biomasses burned by fire events (i.e.,  $L^- + U^- - L^+ - U^+$ ) can often be approximated by three power

187 laws (see Fig. 5 in Maggi and Rinaldi (2006)). This result is somehow  
188 similar to that shown in Fig. 1(b), where, however, the fire intensities  
189 are identified with burned areas.

190 It is worth noticing that in none of the above mentioned studies  
191 the models have been validated against the data collected on a specific  
192 forest site. This is perfectly in line with the aim of the studies, which  
193 was to show that the models could produce fire regimes similar to those  
194 qualitatively observed in various biomes of the world.

195 **Analysis of a spatially extended two-layer forest fire model**

196 The models reviewed in the previous section are definitely poor and  
197 over-simplified from a biological point of view even if they support to a  
198 certain extent some of the characteristics of forest fires emerging from  
199 field data. SOCFF models reduce the growth of vegetation to a sort  
200 of unrealistic ballet of trees born in empty sites and then burned by  
201 lightning, without giving any role to important physical factors such  
202 as quantity of dead biomass on the ground or age of the plants which  
203 are known to control the ignition of a fire and its propagation (Vie-  
204 gas (1998)). By contrast, two-layer models are simply inappropriate  
205 for describing properties concerning burned areas because they do not  
206 explicitly contain space.

207 We therefore focus on a promising mix of the above models by spa-  
208 tially extending on a square lattice with  $L^2$  sites, the two-layer forest-  
209 fire model. Thus, eq. (2) holds at each site, characterized however by  
210 a different standing state  $(L, U)$ , and when the biomasses in one site  
211 reach the pre-fire manifold  $\mathcal{X}^-$ , a fire is ignited in that site of the for-  
212 est and the biomasses of that site are reduced in accordance with the  
213 map described in Fig. 3(a). Moreover, the fire propagates to neigh-  
214 boring sites provided the vegetation in those sites is almost ready to  
215 burn, i.e. provided the biomasses  $(L, U)$  are  $\varepsilon$ -close to the pre-fire man-  
216 ifold  $\mathcal{X}^-$ . In order to simplify the dynamics we assume, in accordance  
217 with Drossel and Schwabl (1992), that the propagation is a sort of in-  
218 stantaneous avalanche, since the time in which a forest cluster burns  
219 down is much shorter than the time in which a tree grows. This means  
220 that when the pre-fire manifold is reached at one site, the fire instan-

221 taneously propagates to an entire forest cluster delimited by sites in  
222 which the biomasses ( $L, U$ ) are at least  $\varepsilon$ -far from the pre-fire manifold  
223  $\mathcal{X}^-$ . Thus, the area burned by fire can be measured by the number of  
224 sites in the cluster.

225 Long simulations of the model allow one to generate long time se-  
226 ries of fires with associated burned areas and times of occurrence. Since  
227 simulations involve time-discretization, it can happen (very rarely how-  
228 ever) that two fires occur at the same time. In these cases one of the  
229 two fires is simply delayed of one time step.

230 In order to avoid finite-size effects we have been forced to work with  
231 large lattices and this is why, in order to keep computational effort un-  
232 der control, we have selected the model described in Maggi and Rinaldi  
233 (2006) which involves  $2L^2$  differential equations, i.e. one half of those  
234 that would be required by the model proposed in Casagrandi and Ri-  
235 naldi (1999). Simulations must be very long because transients toward  
236 attractors of the extended forest model can be extremely long, in par-  
237 ticular when the local dynamics, i.e. the dynamics of a single isolated  
238 site, are chaotic (see Fig. 4 which shows that a reliable estimate of the  
239 mean and standard deviation of the burned areas is obtained only after  
240 three hundred thousand years!).

241 [Figure 4 about here.]

242 Despite these computational difficulties, we have been able to per-  
243 form reliable statistics of the burned areas and of the time of occurrence  
244 of the fires for different values of the parameters of the model. In par-  
245 ticular, we have varied the parameter  $\varepsilon$  that controls the tendency of  
246 the fire to penetrate into parts of the forest which are not yet ready to  
247 burn. Obviously, this parameter depends upon the dominant species  
248 present and can therefore vary remarkably in particular at continental  
249 scale. Higher values of  $\varepsilon$  indicate lower resistance to fire propagation,  
250 i.e. lower tree immunity, as defined in Drossel and Schwabl (1993), Al-  
251 bano (1995) and Song et al. (2001) in their studies on SOCFF models.  
252 Higher values of  $\varepsilon$  should therefore facilitate the occurrence of larger  
253 fires and this is, indeed, what we have systematically found with our  
254 simulations, as shown in Fig. 5 obtained for parameter values in the  
255 range suggested in Maggi and Rinaldi (2006) for Mediterranean forests.

257     Figure 5 shows that the three basic types of distributions identified  
 258     through empirical studies (see Fig. 1) can be produced by our model by  
 259     varying the control parameter  $\varepsilon$ . Another interesting property of our  
 260     model is that large fires, which are associated with the steepest slopes  
 261     of the distributions of the burned areas, are mixed fires, while a relevant  
 262     percentage of the small fires are surface fires. In other words, the results  
 263     suggest that the existence of different slopes in the distributions of the  
 264     burned areas might be due to the existence of differently structured  
 265     fires. This has also been suggested in Schenk et al. (2000) but with  
 266     totally different and less biologically based arguments.

267     Finally, long series of model generated fires have allowed us to esti-  
 268     mate the probability density  $D(\tau|s)$  of the time intervals  $\tau$  separating  
 269     successive fires with burned areas greater than or equal to  $s$ . A typical  
 270     result of this analysis is shown in Fig. 6(a) which compares favourably  
 271     with Fig. 2(a).

273     This means that our model captures also the processes that control the  
 274     times of occurrence of the fires and not only the mechanism regulating  
 275     the severity of the fires, i.e. the burned areas. But the qualitative  
 276     agreement of our model with the empirical evidence goes even further.  
 277     In fact, the similarity of Fig. 2(b) with Fig. 6(b), obtained through  
 278     simulation, proves that the model is endowed with the scaling law (1)  
 279     discovered on purely empirical grounds (Corral et al. (2008)).

280     **Concluding remarks**

281     We have shown that all statistical properties of forest fires discovered  
 282     in the last decade through the analysis of available data can be derived  
 283     from a biological based model in which the three phases of vegetational  
 284     growth, fire ignition and fire propagation are clearly identified. In such  
 285     a model the forest extends over a square lattice of  $L^2$  sites and is com-  
 286     posed of a lower and an upper layer. The two layers grow logistically,  
 287     but the upper one reduces the light available to the lower one, thus  
 288     damaging its growth. Fires are devastating instantaneous events that

289 occur only when the mix of biomasses of the two layers reach partic-  
290 ular values. The rationale for this assumption is as follows. We know  
291 (see, for example, Viegas (1998)) that fire ignition in a forest is possible  
292 only if dead biomass on the ground is above a certain threshold, but  
293 since the biochemical processes regulating the mineralization of dead  
294 biomass are relatively fast with respect to plant growth (Esser et al.  
295 (1982); Seastedt (1988)) it can be reasonably assumed that the rate  
296 of mineralization (proportional to the amount of dead biomass) equals  
297 the inflow rate of new necromass, which, in turn, is proportional to the  
298 standing biomass in the two layers. Thus, in conclusion, the biomasses  
299 of the two layers are appropriate indicators of fuel on the ground, so  
300 that fire ignition is possible only at sites where the standing biomasses  
301 reach specific conditions (called pre-fire conditions). When the fire is  
302 ignited at one site, it immediately propagates to the neighbouring sites  
303 if these are  $\varepsilon$ -close to their pre-fire conditions and this process is re-  
304 peated in an avalanche like manner and stops only when the burning  
305 cluster is delimited by sites which are  $\varepsilon$ -far to their pre-fire conditions.

306 The combination of these slow and fast processes determines the  
307 behavior of the whole forest model which for parameter values in the  
308 ranges suggested in Maggi and Rinaldi (2006) for Mediterranean forests  
309 turns out to be chaotic. In other words, the slow and continuous growth  
310 of the two vegetational layers is punctuated by fires which occur in an  
311 apparently random way in space and time and has statistical properties  
312 consistent with those discovered empirically.

313 It is important to remark that the model proposed in this paper is  
314 nothing but the extension to a network of sites of the minimal model  
315 proposed in Maggi and Rinaldi (2006) for a single site. In other words,  
316 the model is still a minimal model that, as such, can not be calibrated  
317 for performing real time fire predictions in any specific forest, but rather  
318 be used to characterize and classify the fire regimes of large classes of  
319 forests.

320 It is also interesting to remark that the model is fully determin-  
321 istic and spatially homogeneous, so that the emergence of the above  
322 statistical properties does not seem to be necessarily related with the  
323 randomness of meteorological conditions (soil moisture, wind speed, ...)

324 or with geophysical heterogeneity. However, in accordance with Bessie  
325 and Johnson (1995) and Minnich and Chou (1997), we firmly believe  
326 that meteorological randomness and geophysical heterogeneity should  
327 amplify the chaoticity generated by the deterministic mechanisms of  
328 growth, ignition and propagation we have considered. Checking if this  
329 is true could be an interesting point for further investigation, in partic-  
330 ular for assessing the impact of environmental change on fire regimes.  
331 But certainly more interesting would be to try to explain with the model  
332 important regional characteristics of fire regimes that have been discov-  
333 ered from data. For example, the east to west gradient of the slopes of  
334 the power-law distribution across US (Malamud et al. (2005)), might  
335 be a consequence of a similar gradient in some of the parameters of the  
336 model, that control the slopes of the distributions.

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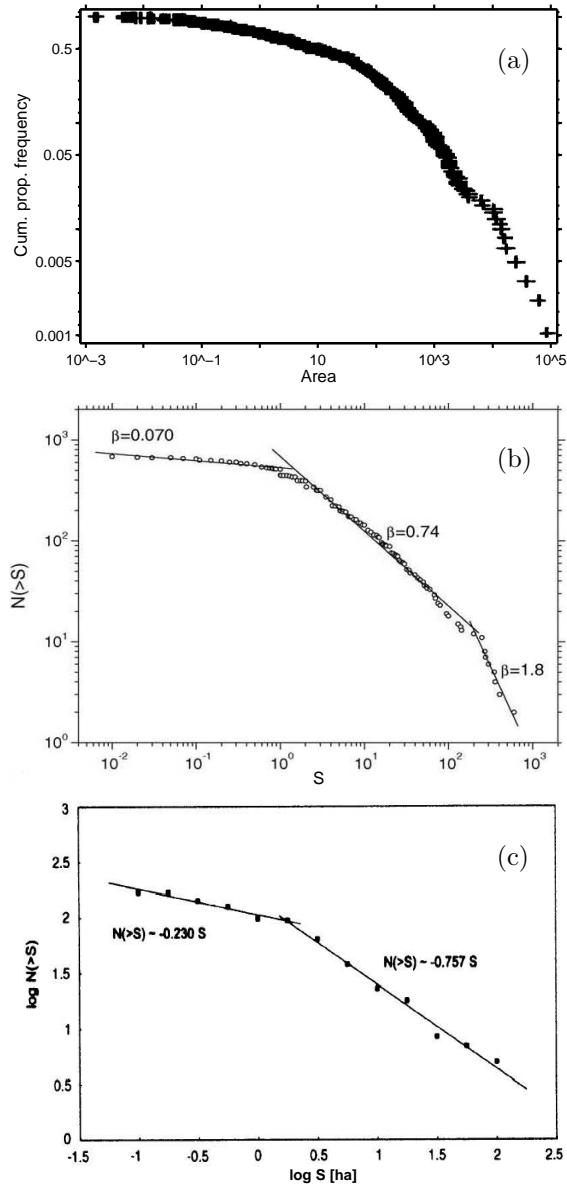


Figure 1: Three examples of cumulative distributions of burned areas obtained from data: (a) Clearwater National Forest (US) redrawn from Reed and McKelvey (2002); (b) Gargano (Italy) redrawn from Telesca et al. (2005); (c) Venaco (Corse, France) redrawn from Ricotta et al. (2001). The distribution in (a) cannot be approximated with a power law, while the distributions in (b) and (c) are approximated with three and two power laws, respectively.

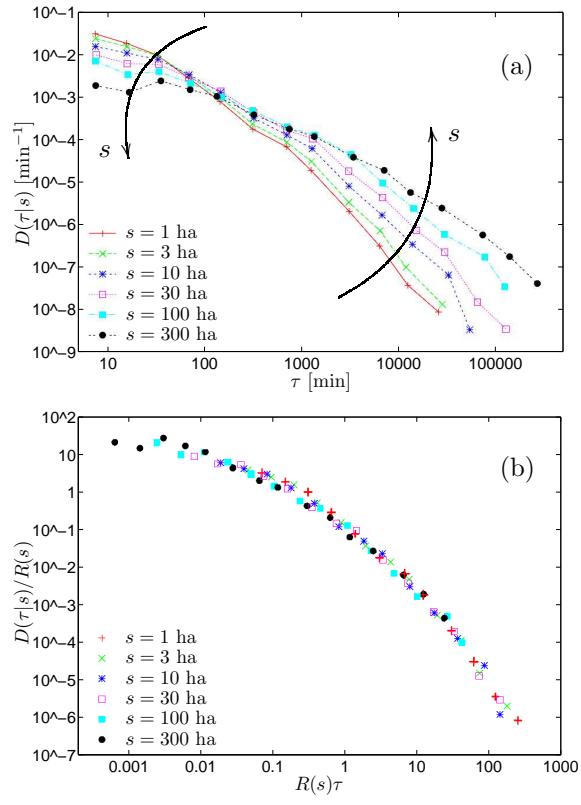


Figure 2: Results of the analysis of the time intervals  $\tau$  separating successive fires in Italy (redrawn from Corral et al. (2008)). (a) Probability densities  $D(\tau|s)$  for different minimum burned areas  $s$ . (b) The previous densities after rescaling by the mean fire rate  $R(s)$  (notice that the rescaling yields dimensionless axes).

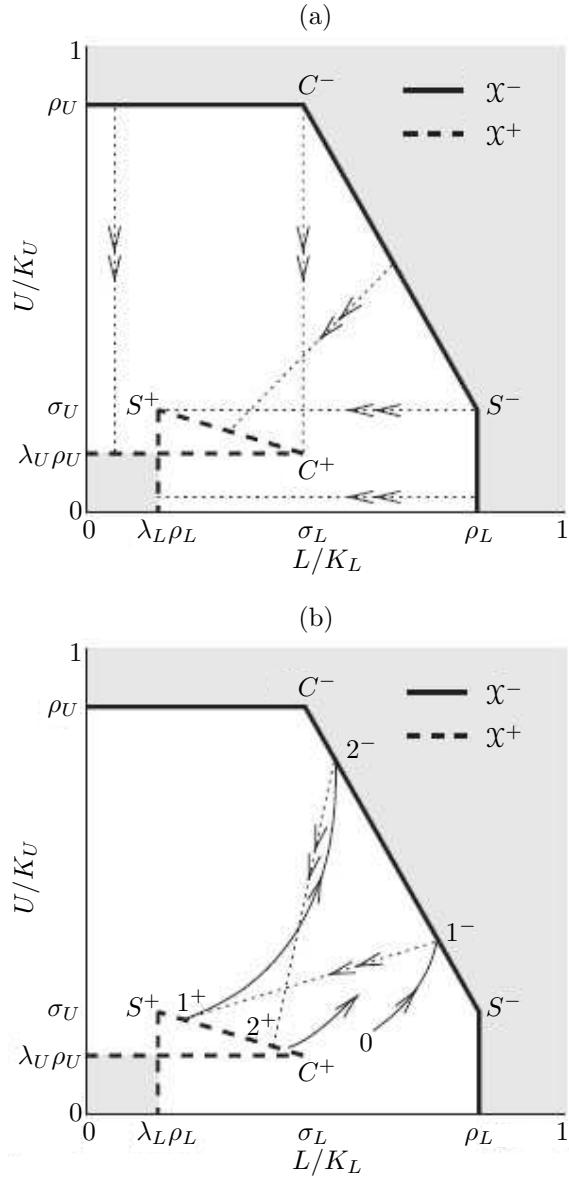


Figure 3: Two-layer model behavior. (a) The pre- and post-fire manifolds  $\mathcal{X}^-$  and  $\mathcal{X}^+$ ; the dotted lines with double arrows are instantaneous transitions from  $\mathcal{X}^-$  to  $\mathcal{X}^+$  due to a fire; horizontal (vertical) lines correspond to surface (crown) fires; oblique lines starting from the segment  $C^- S^-$  of  $\mathcal{X}^-$  correspond to mixed fires. (b) State portrait of the model; continuous lines with a single arrow represent the growing phase of the forest and are described by eq. (2).

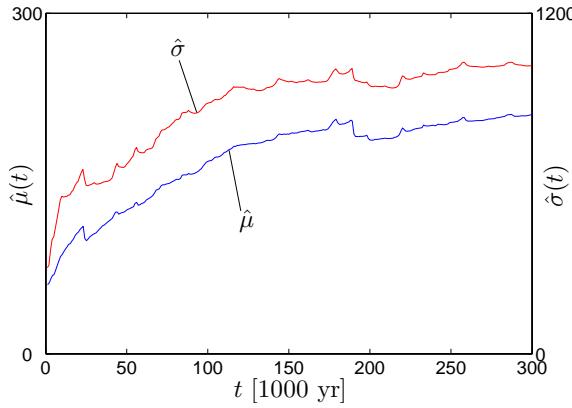


Figure 4: Estimate  $\hat{\mu}$  and  $\hat{\sigma}$  of the mean and standard deviation of the burned areas as a function of the observation time for the model with  $\varepsilon = 0.08, r_L = 3/8, r_U = 1/16, K_L = K_U = 1, \alpha = 129/800, \rho_L = 0.85, \rho_U = 14/15, \sigma_L = 0.6, \sigma_U = 0.35, \lambda_L = \lambda_U = 10^{-4}$ .

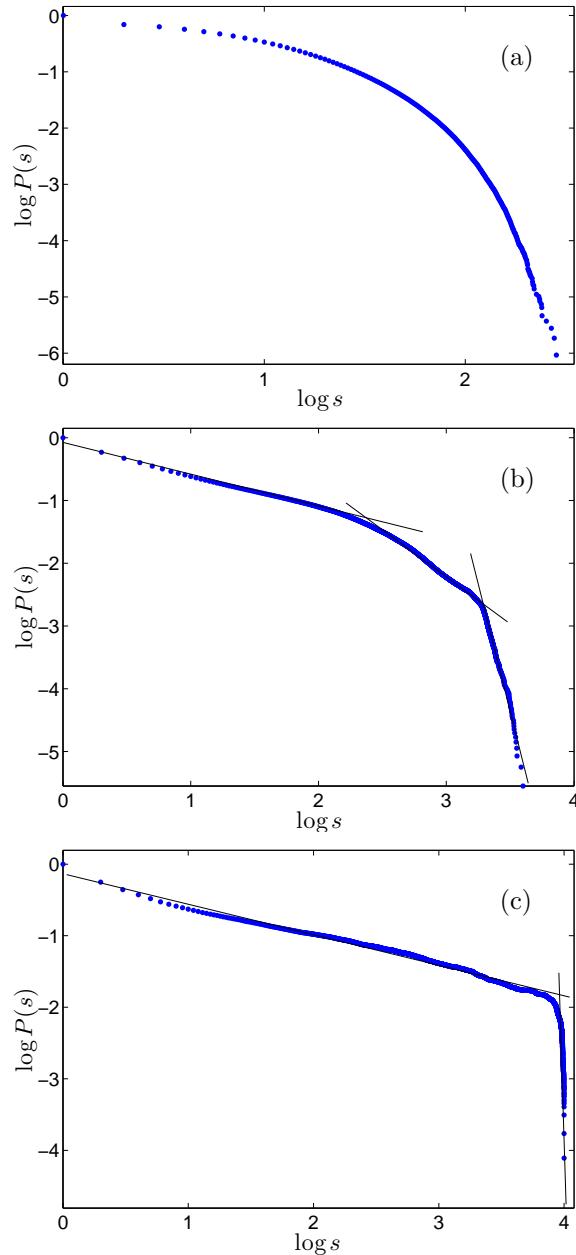


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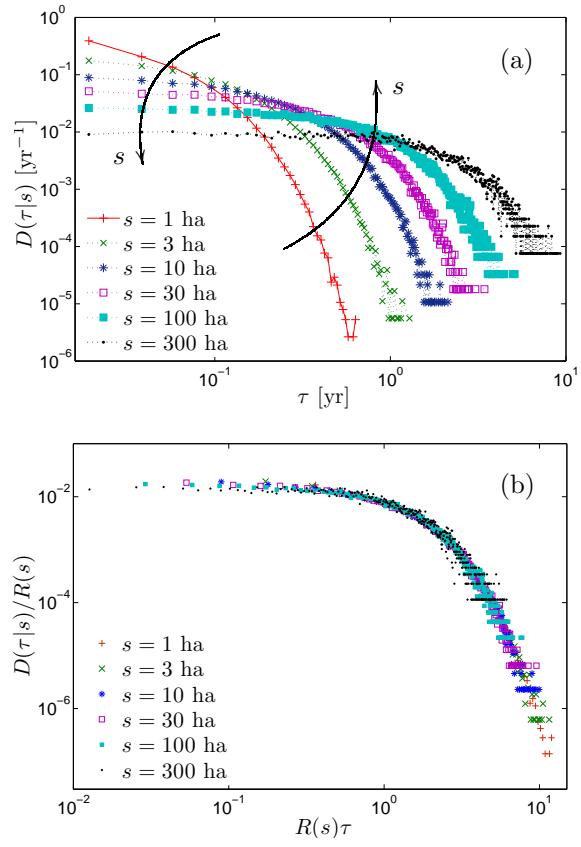


Figure 6: Results of the analysis of the time intervals  $\tau$  separating successive fires generated by the model with parameter values as specified in the caption of Fig. 4. (a) Probability densities  $D(\tau|s)$  for minimum burned areas  $s$ . (b) The previous densities after rescaling by the mean fire rate  $R(s)$ .